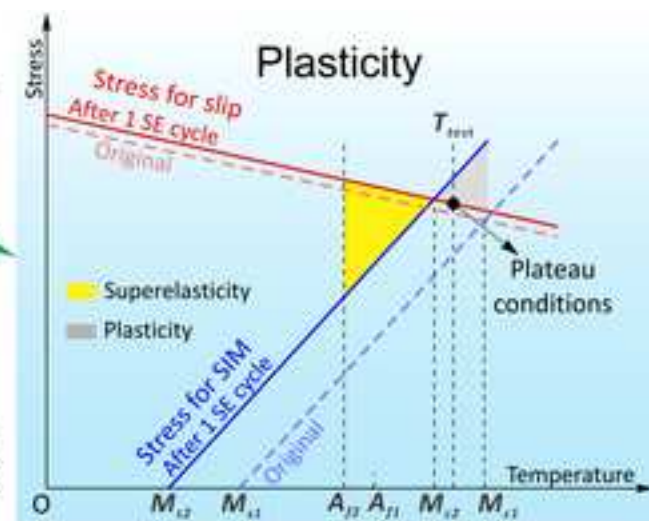
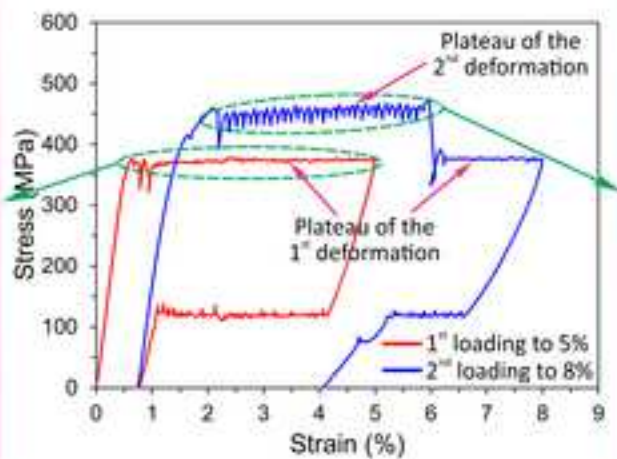
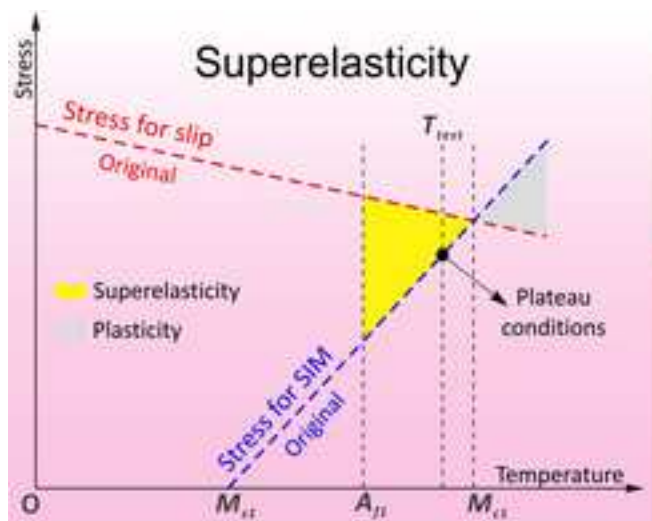


**This item is the archived peer-reviewed author-version of:**

Anomalous stress-strain behavior of NiTi shape memory alloy close to the border of superelastic window

**Reference:**

Wang Xiebin, Yao Xiayang, Schryvers Dominique, Verlinden Bert, Wang Guilong, Zhao Guoqun, Van Humbeeck Jan, Kustov Sergey.- Anomalous stress-strain behavior of NiTi shape memory alloy close to the border of superelastic window  
Scripta materialia - ISSN 1872-8456 - 204(2021), 114135  
Full text (Publisher's DOI): <https://doi.org/10.1016/J.SCRIPTAMAT.2021.114135>  
To cite this reference: <https://hdl.handle.net/10067/1816580151162165141>



## Anomalous stress-strain behavior of NiTi shape memory alloy close to the border of superelastic window

Xiebin Wang<sup>a,b,\*</sup>, Xiayang Yao<sup>c</sup>, Dominique Schryvers<sup>c</sup>, Bert Verlinden<sup>d</sup>, Guilong Wang<sup>a</sup>, Guoqun Zhao<sup>a</sup>, Jan Van Humbeeck<sup>d</sup>, Sergey Kustov<sup>e</sup>

<sup>a</sup>Key Laboratory for Liquid-Solid Structural Evolution and Processing of Materials (Ministry of Education), Shandong University, Jingshi Road 17923, Jinan 250061, China.

<sup>b</sup>Suzhou Institute of Shandong University, Ruoshui Road 388, Suzhou 215123, China.

<sup>c</sup>Electron Microscopy for Materials Science (EMAT), University of Antwerp, Groenenborgerlaan 171, Antwerp B-2020, Belgium.

<sup>d</sup>Department of Materials Engineering, University of Leuven (KU Leuven), Kasteelpark Arenberg 44 bus 2450, Heverlee B-3001, Belgium.

<sup>e</sup>Departament de Física, Universitat de les Illes Balears, Cra Valldemossa km 7.5, Palma de Mallorca E07122, Spain.

\*Corresponding author: xiebin.wang@email.sdu.edu.cn, wangxiebin@sdu.edu.cn (X. Wang)

### Abstract

In this work, we report an anomalous phenomenon on superelastic cycling of NiTi shape memory alloys when deforming at the temperature close to the border of superelastic window. New unexpected effects are found: (i) critical stress for inducing martensite transformation during the second loading cycle is higher than that of the first cycle; (ii) the plateau stress of the second cycle decreases to the original level when the strain overcomes the limit of the first cycle; (iii) transition from good superelasticity in the first cycle to fully irreversible strain in the second. We propose that defects generated during the first superelastic cycle close to the border of superelastic window impede following stress-induced martensitic transformations, leading to the increase of critical stress beyond yield stress of the B2 matrix, and thus functional fatigue of NiTi alloys.

**Keywords:** Shape memory alloy, NiTi, Plastic deformation, Superelasticity, Functional fatigue.

1 Near equiatomic NiTi shape memory alloy transforms thermoelastically between a B2  
2 structured austenite (A) and B19' martensite (M), giving rise to the effects of shape memory  
3 and superelasticity [1], which are widely employed in practical applications [2,3]. Apart from  
4 the unique functional properties, the mechanism of plastic deformation in NiTi alloys has  
5 attracted increasing interest [4-10], since plastic activities accompanying the martensitic  
6 transformation (MT) are the main reason for the functional degradation of NiTi alloys [10-14].  
7 Thus, understanding the origin of these plastic activities is a key step to improve the functional  
8 stability of NiTi alloys.  
9

10  
11  
12  
13  
14  
15  
16 The origin of the plastic activities accompanying the MT in NiTi has been studied both through  
17 experimental work and simulation [4-6]. Because the B2 and B19' phases have different lattice  
18 structure [1], the strain compatibility of the coherent A/M interface during transformation  
19 has been addressed. It is suggested that, in order to accommodate the local strain gradient,  
20 plastic slip during the forward A→M transformation may occur either at the A/M interface  
21 [7,15,16], or in the austenite [10,12,16-21] / martensite [9,22] matrix adjacent to the A/M  
22 interface. The dislocation slip occurring during the reverse transformation is also discussed  
23 [16-18,23]. Deformation twinning is another possible contribution of the irreversible strain  
24 [8,24-28], especially when deforming at elevated temperatures [8] or subjected to high  
25 loading cycles [25].  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35

36  
37 However, to the best of our knowledge, the origin of plastic activities in NiTi alloys during  
38 stress-induced MT is still ambiguous, mainly due to the complexity of this process, involving  
39 phase transformation, dislocation slip, martensite twin migration, and deformation twinning.  
40 Moreover, the deformation behavior of NiTi is sensitive to many other factors, for instance,  
41 crystallographic orientation and texture [29-31], presence of inclusions or precipitates [32-  
42 34], grain size [35,36], thermo-mechanical history (e.g., post-deformation annealing, aging)  
43 [37,38], load path [39,40], strain rate [41-43], strain amplitude [8], and testing temperature  
44 [5,23]. The experimental work with unconventional deformation behavior might help to  
45 understand the plastic deformation in NiTi alloys, as previously suggested by Chen *et al.* [8].  
46  
47  
48  
49  
50  
51  
52  
53  
54

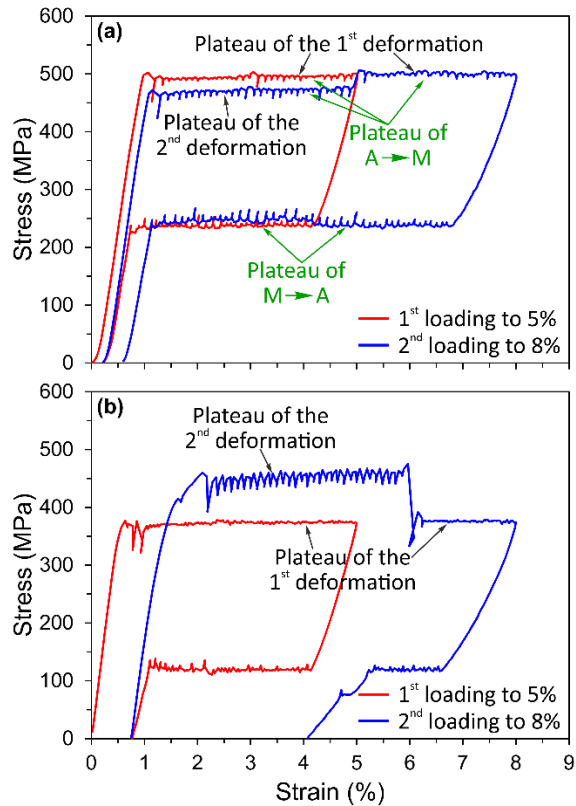
55 In the present work, following this idea, we study superelastic behavior of NiTi upon  
56 approaching the border of superelastic window in order to study the transition from  
57 superelastic to plastic deformation. We show that the role of defects created during stress-  
58  
59  
60  
61  
62  
63  
64  
65

1 induced MT changes close to this border, and the effect should be taken into account in  
2 functional fatigue of NiTi alloys.  
3

4  
5 A cold drawn Ti-50.8 at.% Ni wire of 100  $\mu\text{m}$  diameter (Memory-Metalle GmbH, Germany)  
6 was used. The as-received wire was annealed for 30 min at different temperatures (773, 823  
7 and 873 K) under protective Ar atmosphere and water-quenched at room temperature. Major  
8 attention was paid to 873 K-annealed samples, while the 773 K and 823 K-annealed ones  
9 served as a reference. The 873 K-annealed sample shows a direct MT start temperature ( $M_s$ )  
10 at 227 K and reverse MT finish temperature ( $A_f$ ) at 257 K (Fig. S1 in supplement). Mechanical  
11 tests were carried out using a TA Q800 dynamic mechanical analyzer under a constant strain  
12 rate of  $1.67 \times 10^{-4} \text{ s}^{-1}$ . Electrical resistance (ER) measurements were performed to characterize  
13 the transformation behavior. The resistivity was registered using AC excitation and lock-in  
14 detection in a home-made four-probe measuring system under a cooling/heating rate of 2 K  
15  $\text{min}^{-1}$ . The microstructure was investigated by means of electron backscatter diffraction (EBSD)  
16 and transmission electron microscopy (TEM). EBSD tests were performed using a TSL/EDAX  
17 system on a FEI Nova NanoSEM 450 microscope. For EBSD observations, the surface parallel  
18 to the wire axis was prepared using a JEOL IB-09010CP cross-section polisher. TEM was  
19 conducted using a FEI Tecnai G2 microscope operated at 200 kV. The TEM foils, taken along  
20 the wire axis, were prepared by focused ion beam (FIB) with the 'lift-out' procedure in a FEI  
21 Helios 650 NanoLab SEM/FIB instrument, an example of the TEM sample is shown in Fig. S2  
22 in supplement.  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40

41 Figure 1a shows typical reference superelastic stress-strain curves in two consecutive cycles  
42 for an 823 K-annealed sample. As the stress reaches a critical value, the stress-induced  
43 martensite transformation (SIMT) occurs through a Lüders-like mechanism [44,45], giving rise  
44 to a stress plateau during loading. Except for a small amount of residual strain, most of the  
45 strain is recovered during unloading due to the reverse  $M \rightarrow A$  transformation, which also  
46 produces a stress plateau. The residual strain is mainly caused by the plastic activities  
47 associated with the forward and/or reverse transformation [4-6]. These plastic activities  
48 induce an internal stress field, which can assist the formation of martensite variants [15,19,46]  
49 in consecutive SIMTs. As a result, the critical stress of SIMT is lower in the second cycle than  
50 in the first cycle (blue curve in Fig. 1a), but only for the strain below the maximum reached in  
51 the first cycle. The critical stress **increases to the original level** as soon as the strain exceeds  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65

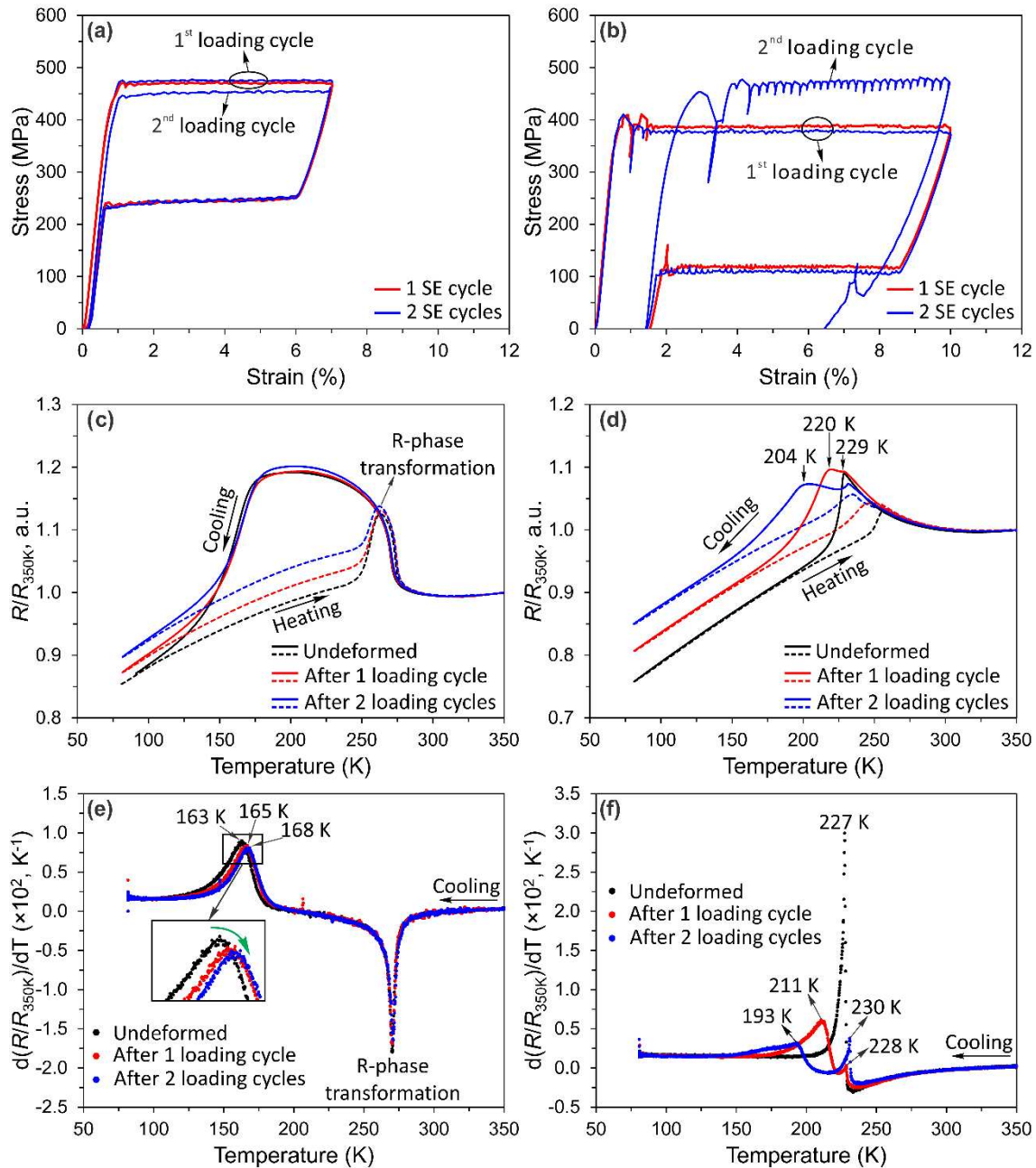
the limit of the first cycle. A relatively small increase of irrecoverable strain (around 0.5%) is observed after the second cycle.



**Figure 1.** Stress-strain curves of a cold-drawn Ti-50.8 at.% wire after annealing for 30 min at (a) 823 K, and (b) 873 K. The samples are loaded at 298 K to a strain of 5% and 8% during the first and second cycle, respectively.

Figure 1b shows the cyclic deformation behavior of the 873 K-annealed sample. The results exemplify the unusual stress-strain behavior at 298 K, and the striking observations are: (i) the plateau stress of the second cycle is higher than that of the first cycle, which is different from the conventional superelastic behavior (Fig. 1a); (ii) opposite to the behavior of the 823 K-annealed sample (Fig. 1a), the plateau stress **drops to the original level** of the first cycle, when the strain exceeds the limit of the first cycle. Thus, in the second loading cycle, in violation of the expected natural sequence, the deformation proceeds first at a high stress and then at a substantially lower one (Fig. 1b blue curve, and Fig. S3 in supplement); (iii) the first cycle shows good superelasticity, but the second one shows recoverable strain comparable only with the strain added in the second cycle (Fig. 1b and Fig. S4 in supplement). The strain around 4.3% that was essentially reversible (superelasticity) in the first cycle

1 becomes irreversible in second cycle. Incremental loading-unloading cycles up to a strain of  
 2 25% (Fig. S5 in supplement), indicate that these unusual features persist in consecutive cycles  
 3 until the strain becomes purely plastic.  
 4  
 5  
 6



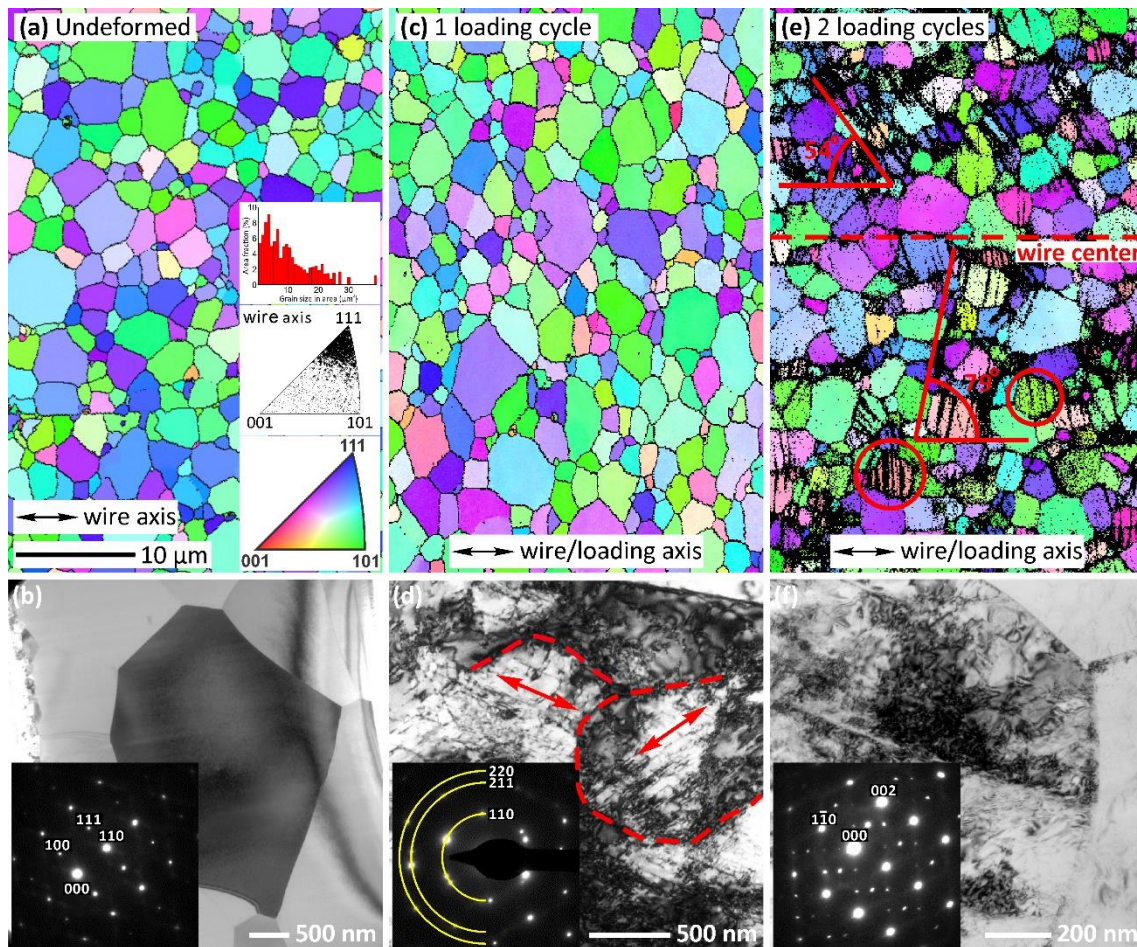
**Figure 2.** Stress-strain curves of a cold drawn Ti-50.8 at.% Ni wire after annealing for 30 min at (a) 773 K and (b) 873 K. Two samples were subjected to tensile loading-unloading tests at 298 K for 1 and 2 cycles, respectively. The electric resistance curves of (c) the 773K-annealed and (d) the 873K-annealed samples before and after one and two loading cycles. (e) and (f) show temperature derivative of the electrical resistance from (c) and (d), respectively.

1 Figure 2 shows the preparatory stress-strain cycles for cold drawn Ti-50.8 at.% Ni wire after  
2 annealing at 773 (a) and 873 K (b). The sample annealed at 773 K shows a conventional  
3 superelastic response, as frequently reported in literature [1,47], while the one annealed at  
4 873 K shows the abovementioned counterintuitive behavior. The influence of superelastic  
5 loading on the temperature-induced MT was characterized by ER measurements. According  
6 to Ref. [48],  $dR/dT$  perfectly correlates with the transformation rate in NiTi alloys. Resistivity  
7 measurements detect at around 270 K well documented the R-phase transformation in the  
8 773 K-annealed sample (Fig. 2c), which is due to the incomplete annealing and formation of  
9 coherent  $Ni_4Ti_3$  precipitates during annealing below 873 K [1,47,49]. On further cooling the R-  
10 phase transforms into martensite at around 160 K. The R-phase to martensite transformation  
11 shifts gradually to higher temperatures upon superelastic cycling (Fig. 2e). The  $dR/dT$  peak  
12 temperature is 163 K for the as-annealed state, shifting to 165 K and to 168 K after one and  
13 two superelastic cycles, respectively. The increase of transformation temperature is  
14 consistent with the decrease of the plateau stress, according to the Clausius-Clapeyron type  
15 relation [1]. It has frequently been suggested that a “directional” internal stress field could be  
16 generated after superelastic loading, due to the plastic activities associated with SIMT  
17 [15,19,30,46]. The stress field can promote the MT and thus leads to the decrease of the  
18 plateau stress (Fig. 2a) and increase of the transformation temperatures (Fig. 2c and 2e) after  
19 superelastic cycling.  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36

37  
38 Figures 2d and 2f show the  $R(T)$  and its temperature derivative for the 873 K-annealed sample,  
39 also in as annealed state and after 1 and 2 superelastic cycles (Fig. 2b). The R-phase  
40 transformation is absent, due to the full annealing of the cold-worked microstructure at 873  
41 K [37,50,51]. A remarkable observation is the strong splitting of the direct MT into two stages  
42 in deformed samples. The temperature of the first step in deformed samples is the same as  
43 in the as-annealed state. Therefore, the two steps of the MT in deformed samples should be  
44 associated with parts of the material that suffered the SIMT and remained not transformed.  
45 Crucial observation is that the second step of the MT, related with the part of the sample that  
46 suffered SIMT, is strongly shifted to lower temperatures by 16 and 34 K after one and two  
47 loading cycles, respectively (Fig. 2f). The decrease of the MT temperature is also revealed by  
48 DSC measurement (Fig. S6 in supplement). Thus, the effect of SIMT during cyclic loading on  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65



temperature-induced MT is opposite in samples annealed at 773 and 873 K, pointing to the different status of defects generated during superelastic cycling.



**Figure 3.** (a) EBSD orientation map and (b) TEM bright field image of a cold drawn Ti-50.8 at.% Ni wire after annealing at 873 K for 30 min. The grain size distribution chart, inverse pole figure with respect to the wire axis, and the color code are inserted in (a). (c) EBSD map and (d) TEM bright field image of the as-annealed sample after 1 loading-unloading cycle at 298 K. (e) EBSD map and (f) TEM bright field image of the as-annealed sample after 2 loading-unloading cycles at 298 K. The dash lines in (d) roughly indicate some grain boundaries.

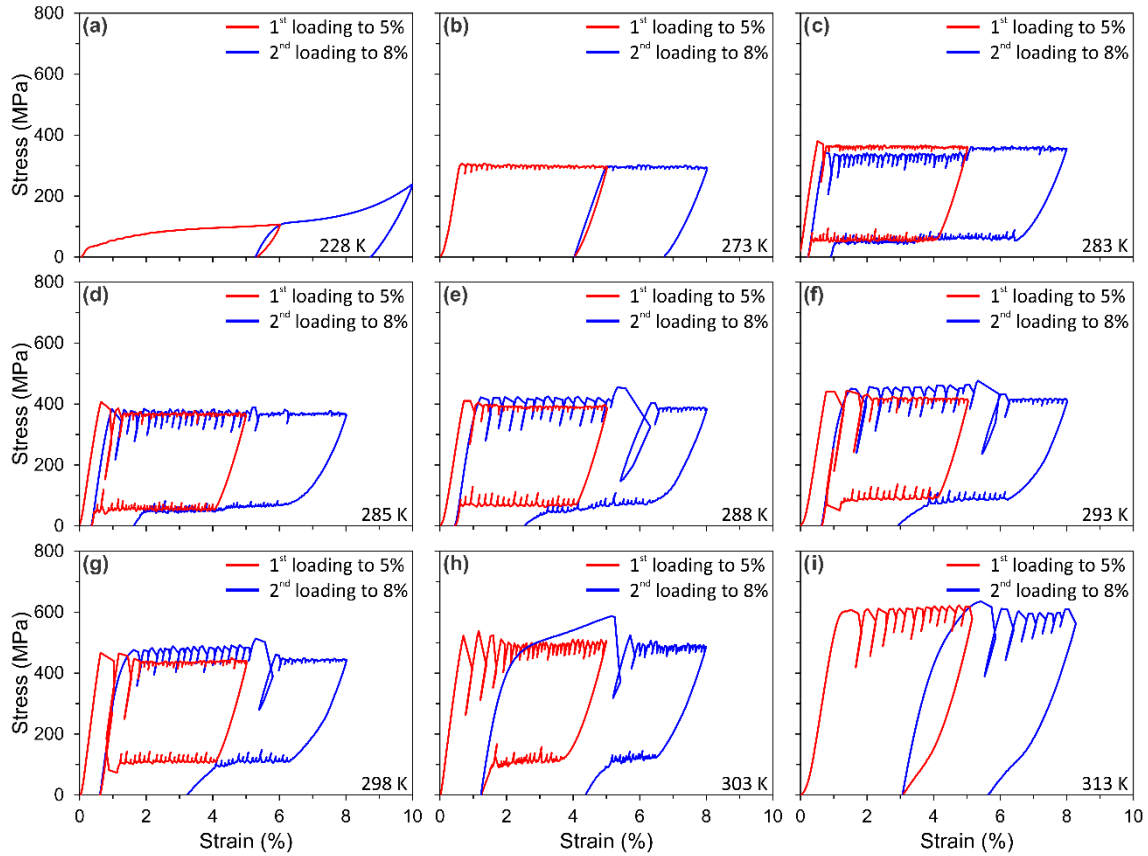
Figure 3 shows the microstructure of the 873 K-annealed sample, which possesses the uncovered anomalous stress-strain behavior (Fig. 1b). The EBSD map of as-annealed sample in Fig. 3a shows a fully recrystallized microstructure with equiaxed grains. The average grain size was manually determined as 2.0  $\mu\text{m}$  using the line intercept method described by Sutou *et al.* [52]. Figure 3a also gives the inverse pole figure with reference to the wire axis, and a strong  $\langle 111 \rangle_{B2}$  fiber texture is revealed. TEM observations (Fig. 3b) reveal very low dislocation

1 density and sharp grain boundaries. The corresponding selected area diffraction (SAD)  
2 pattern in Fig. 3b confirms that the as-annealed sample is in B2 austenite phase at room  
3 temperature.  
4

5  
6 EBSD results for the sample after 1 superelastic cycle (Fig. 3c) do not show any obvious  
7 microstructural difference as compared with the as-annealed state (Fig. 3a). Nevertheless, a  
8 significant increase of dislocation density is revealed by TEM (Fig. 3d). The massive dislocation  
9 slip results in an irrecoverable strain of 1.4% after unloading (Fig. S7 in supplement). The  
10 dislocation network consists mainly of arrays, frequently reported after superelastic cycling  
11 [16-18,20]. The direction of dislocation arrays, as marked by the arrows in Fig. 3d, varies in  
12 different grains, confirming a strong influence of crystallographic orientation [16-18,20,53].  
13  
14  
15  
16  
17  
18  
19  
20

21 After 2 superelastic cycles, EBSD reveals massive traces of plastic deformation (Fig. 3e). The  
22 deformation bands are inclined clockwise and counterclockwise with respect to the wire axis  
23 below and above the wire axis, respectively. This feature points to the formation of a cone-  
24 shaped A/M interface during SIMT in the wire, as reported by Sedmark *et al.* [54]. Unlike the  
25 frequently reported inclination angle around 55° for the macroscopic Lüders bands [21,55],  
26 the inclination angle of the deformation bands in Fig. 3e varies largely over different grains.  
27 TEM observations (Fig. 3f) reveal a high density of entangled dislocations. Straight dislocation  
28 segments are not observed, in contrast to the sample subjected to 1 superelastic cycle.  
29  
30  
31  
32  
33  
34  
35  
36  
37

38 Figure 4 shows the cyclic stress-strain response of 873 K-annealed samples tested at different  
39 temperatures. It indicates that the appearance of the counterintuitive phenomenon depends  
40 highly on the testing temperature. The superelasticity is absent when testing at 228 and 273  
41 K, because the low testing temperature prohibits the reverse transformation. Superelasticity  
42 is observed when testing at and above 283 K. The plateau stress of both the first and second  
43 loading cycle increases with the increase of testing temperature. However, the plateau stress  
44 of the second cycle increases faster than that of the first cycle. The sample tested at 283 K  
45 shows a conventional superelastic response (Fig. 4c). At 285 K the two loading cycles show  
46 similar plateau stress (Fig. 4d). The uncovered unusual behavior is observed for tests between  
47 288 and 303 K (Fig. 4e-4g). The sample loses completely superelasticity and undergoes plastic  
48 deformation at 313 K (Fig. 4i). Thus, this temperature is considered as the upper limit of the  
49 superelastic window for 873 K-annealed samples.  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65

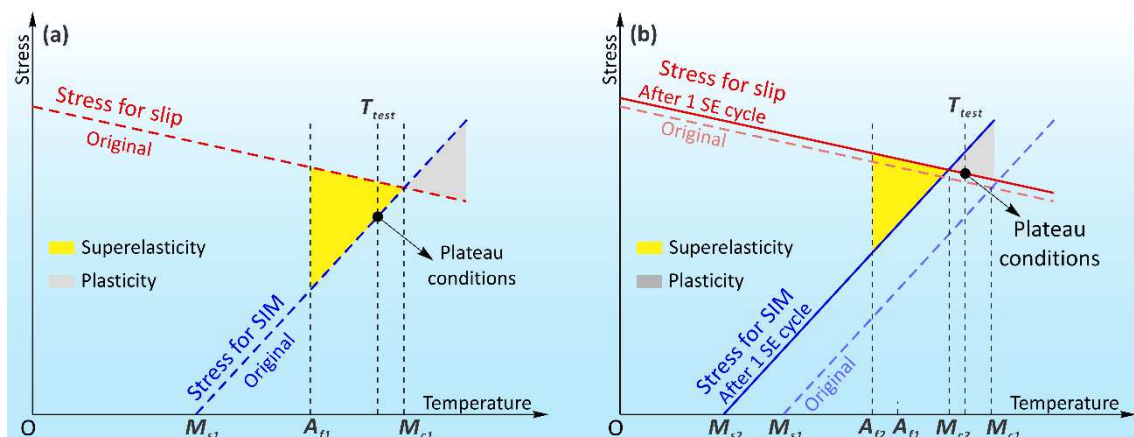


**Figure 4.** The stress-strain curves, tested at different temperatures, of a cold drawn Ti-50.8 at.% Ni wire after annealing at 873 K for 30 min.

Figure 5 explains a possible scenario that explains the occurrence of the anomalous stress-strain behavior just below the high-temperature border of the superelastic window (288-303 K in this work). Figure 5a illustrates the states of the material before the first loading cycle. With the increase of temperature, the critical stress for SIMT increases according to the Clausius-Clapeyron type relation, while the yield strength of the alloy decreases. Hence, we could define a critical temperature  $M_{c1}$ , at which the yield stress equals the stress for SIMT. The sample will lose superelasticity above  $M_{c1}$ , since plastic deformation will occur at a lower stress than SIMT. According to the results in Fig. 4, the testing temperature of 298 K is below but rather close to the  $M_{c1}$  temperature, since the sample loses superelasticity between 303 and 313 K.

The defects generated during the first superelastic cycle alter the state of the material, provoking the decrease of MT temperature (Fig. 2 and Fig. S5) from  $M_{s1}$  to  $M_{s2}$ . Therefore, the  $M_c$  temperature decreases to a lower value,  $M_{c2}$ , as is illustrated in Fig. 5b. Moreover, the

critical stress for SIMT at fixed temperature of 298 K increases for the second loading cycle. Since the testing temperature is close to the high-temperature border of the superelastic window, the critical stress becomes higher than the yield stress of the austenite during the second cycle. That is, the testing temperature becomes higher than the modified  $M_{c2}$ . Therefore, the part of the sample that underwent SIMT in the first cycle undergoes plastic deformation during the second cycle. This interpretation leaves open the most intriguing question: why the SIMT of the “virgin” part of the sample occurs at lower stress *after* the plastic deformation of the transformed part during the second loading cycle (blue curve in Fig. 1b, and Fig. S3 in supplement)? Apparently, the simple interpretation based on Fig. 5 is not able to account for this counterintuitive behavior. Possible research line for the future work is to study the transformation interface during the first loading cycle, which might be the key to understand the uncovered effects. It is also worthy to check whether the reported effects are generic and inherent in other NiTi alloys (e.g. aged or partially annealed samples) at the temperatures close to the corresponding limit of superelastic window.



**Figure 5.** Schematic illustration of the states of the sample (a) before the first loading cycle, and (b) after the first loading cycle.  $M_c$  indicates the critical temperature, at which the critical stress for inducing martensite transformation (SIM) equals the yield stress of the matrix.

In summary, we argue that upon approaching the border of superelastic window, the role of defects created during stress-induced MT changes. Well below the limit of this window, the defects assist consecutive transformations through their stress fields. For stress-induced MT close to the limit of superelastic window, the created defects strongly impede consecutive MT. The change of the role of plastic activities upon approaching the border of superelastic

1 window results in a transition from conventional superelasticity to the unusual stress-strain  
2 behavior with certain counterintuitive features. The different role of SIMT-induced defects,  
3 depending on the distance from the border of superelastic window, might be important in  
4 understanding and controlling functional fatigue of NiTi alloys.  
5  
6

### 7 8 9 **Declaration of interests**

10  
11 The authors declare that they have no known competing financial interests or personal  
12 relationships that could have appeared to influence the work reported in this paper.  
13  
14

### 15 16 17 **Acknowledgements**

18  
19 This work was supported by National Key R&D Program of China [grant number  
20 2018YFB1105100]; National Nature Science Foundation of China [grant number 51905310];  
21 Natural Science Foundation of Shandong Province [ZR2020YQ39]; Natural Science Foundation  
22 of Jiangsu Province [grant number BK20180231]; Key Research and Development Program of  
23 Shandong Province [grant number 2019GGX104065]; the Young Scholars Program of  
24 Shandong University [grant number 2018WLJH24]; and the Research Foundation Flanders  
25 [grant number G.0366.15N]. S. Kustov acknowledges the support from Spanish MCIU-AEI and  
26 FEDER, UE [project RTI 2018-094683-B-C51].  
27  
28  
29  
30  
31  
32  
33  
34  
35

### 36 37 38 **References**

- 39 [1] K. Otsuka, X. Ren, Prog. Mater. Sci. 50 (2005) 511-678.  
40  
41 [2] C. Constant, S. Nichols, É. Wagnac, Y. Petit, A. Desrochers, V. Brailovski, Materialia 9 (2020)  
42 100567.  
43  
44 [3] D. Kapoor, Johnson Matthey Technol. Rev. 61 (2017) 66-76.  
45  
46 [4] P. Sittner, P. Sedlak, H. Seiner, P. Sedmak, J. Pilch, R. Delville, L. Heller, L. Kaderavek, Prog. Mater.  
47 Sci. 98 (2018) 249-298.  
48  
49 [5] P. Chowdhury, H. Sehitoglu, Prog. Mater. Sci. 85 (2017) 1-42.  
50  
51 [6] P. Chowdhury, H. Sehitoglu, Prog. Mater. Sci. 88 (2017) 49-88.  
52  
53 [7] H. Sehitoglu, Y. Wu, S. Alkan, E. Ertekin, Philos. Mag. Lett. 2017 (2017) 1-12.  
54  
55 [8] Y. Chen, O. Molnárová, O. Tyc, L. Kadeřávek, L. Heller, P. Šittner, Acta Mater. 180 (2019) 243-259.  
56  
57 [9] A.S.K. Mohammed, H. Sehitoglu, Acta Mater. 186 (2020) 50-67.  
58  
59 [10] H.M. Paranjape, M.L. Bowers, M.J. Mills, P.M. Anderson, Acta Mater. 132 (2017) 444-454.  
60  
61 [11] Y. Gao, Materialia 6 (2019).  
62  
63  
64  
65

- 1  
2  
3  
4  
5  
6  
7  
8  
9  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65
- [12] L. Heller, H. Seiner, P. Sittner, P. Sedlak, O. Tyc, L. Kaderavek, *Int. J. Plast.* 111 (2018) 53-71.
- [13] Y. Gao, L. Casalena, M.L. Bowers, R.D. Noebe, M.J. Mills, Y. Wang, *Acta Mater.* 126 (2017) 389-400.
- [14] W.-S. Ko, W.S. Choi, G. Xu, P.-P. Choi, Y. Ikeda, B. Grabowski, *Acta Mater.* 202 (2021) 331-349.
- [15] P. Sedmak, P. Sittner, J. Pilch, C. Curfs, *Acta Mater.* 94 (2015) 257-270.
- [16] T. Simon, A. Kröger, C. Somsen, A. Dlouhy, G. Eggeler, *Acta Mater.* 58 (2010) 1850-1860.
- [17] D.M. Norfleet, P.M. Sarosi, S. Manchiraju, M.F.-X. Wagner, M.D. Uchic, P.M. Anderson, M.J. Mills, *Acta Mater.* 57 (2009) 3549-3561.
- [18] M.L. Bowers, X. Chen, M. De Graef, P.M. Anderson, M.J. Mills, *Scripta Mater.* 78-79 (2014) 69-72.
- [19] R.F. Hamilton, H. Sehitoglu, Y. Chumlyakov, H.J. Maier, *Acta Mater.* 52(11) (2004) 3383-3402.
- [20] R. Delville, B. Malard, J. Pilch, P. Sittner, D. Schryvers, *Acta Mater.* 58 (2010) 4503-4515.
- [21] H. Kato, K. Sasaki, *Int. J. Plast.* 50 (2013) 37-48.
- [22] C. Yu, G. Kang, D. Song, Q. Kan, *Int. J. Plast.* 67 (2015) 69-101.
- [23] L. Heller, P. Šittner, P. Sedlák, H. Seiner, O. Tyc, L. Kadeřávek, P. Sedmák, M. Vronka, *Int. J. Plast.* 116 (2019) 232-264.
- [24] O. Benafan, R.D. Noebe, S.A. Padula, A. Garg, B. Clausen, S. Vogel, R. Vaidyanathan, *Int. J. Plast.* 51 (2013) 103-121.
- [25] R. Delville, B. Malard, J. Pilch, P. Sittner, D. Schryvers, *Int. J. Plast.* 27 (2011) 282-297.
- [26] T. Ezaz, H. Sehitoglu, H.J. Maier, *Acta Mater.* 60 (2012) 339-348.
- [27] S. Ii, K. Yamauchi, Y. Maruhashi, M. Nishida, *Scripta Mater.* 49 (2003) 723-727.
- [28] Y. Zhong, K. Gall, T. Zhu, *Acta Mater.* 60 (2012) 6301-6311.
- [29] Z. Wang, J. Chen, C. Besnard, L. Kunčická, R. Kocich, A.M. Korsunsky, *Acta Mater.* 202 (2021) 135-148.
- [30] A.W. Richards, R.A. Lebensohn, K. Bhattacharya, *Acta Mater.* 61 (2013) 4384-4397.
- [31] C.P. Frick, B.G. Clark, A.S. Schneider, R. Maab, S. Van Petegem, H. Van Swygenhoven, *Scripta Mater.* 62 (2010) 492-495.
- [32] P. Chowdhury, L. Patriarca, G. Ren, H. Sehitoglu, *Int. J. Plast.* 81 (2016) 152-167.
- [33] M. Rahim, J. Frenzel, M. Frotscher, J. Pfetzing-Micklich, R. Steegmüller, M. Wohlschlägel, H. Mughrabi, G. Eggeler, *Acta Mater.* 61 (2013) 3667-3686.
- [34] K. Gall, H.J. Maier, *Acta Mater.* 50 (2002) 4643-4657.
- [35] A. Ahadi, Q. Sun, *Acta Mater.* 90 (2015) 272-281.
- [36] J. Chen, H. Yin, Q. Sun, *Acta Mater.* 195 (2020) 141-150.
- [37] X. Wang, S. Kustov, K. Li, D. Schryvers, B. Verlinden, J. Van Humbeeck, *Acta Mater.* 82 (2015) 224-233.
- [38] Y. Kaynak, H. Tobe, R.D. Noebe, H.E. Karaca, I.S. Jawahir, *Scripta Mater.* 74 (2014) 60-63.
- [39] W.-N. Hsu, E. Polatidis, M. Smid, N. Casati, S. Van Petegem, H. Van Swygenhoven, *Acta Mater.* 144 (2018) 874-883.



- 1 [40] J. Arghavani, F. Auricchio, R. Naghdabadi, A. Reali, S. Sohrabpour, *Int. J. Plast.* 26 (2010) 976-  
2 991.
- 3 [41] A. Ahadi, Q. Sun, *Acta Mater.* 76 (2014) 186-197.
- 4 [42] R.R. Adharapurapu, F. Jiang, K.S. Vecchio, G.T. Gray III, *Acta Mater.* 54 (2006) 4609-4620.
- 5 [43] S. Nemat-Nasser, J.Y. Choi, *Acta Mater.* 53 (2005) 449-454.
- 6 [44] G. Tan, Y. Liu, P. Sittner, M. Saunders, *Scripta Mater.* 50 (2004) 193-198.
- 7 [45] J.A. Shaw, S. Kyriakides, *Acta Mater.* 45 (1997) 683-700.
- 8 [46] P.K. Kumar, D.C. Lagoudas, *Acta Mater.* 58 (2010) 1618-1628.
- 9 [47] X. Huang, Y. Liu, *Scripta Mater.* 45 (2001) 153-160.
- 10 [48] S. Kustov, D. Salas, E. Cesari, R. Santamarta, J. Van Humbeeck, *Acta Mater.* 60(6-7) (2012) 2578-  
11 2592.
- 12 [49] Y. Song, M. Jin, X. Han, X. Wang, P. Chen, X. Jin, *Acta Mater.* 205 (2021) 116541.
- 13 [50] X. Wang, C. Li, B. Verlinden, J. Van Humbeeck, *Scripta Mater.* 69 (2013) 545-548.
- 14 [51] X. Wang, Z. Pu, Q. Yang, S. Huang, Z. Wang, S. Kustov, J. Van Humbeeck, *Scripta Mater.* 163  
15 (2019) 57-61.
- 16 [52] Y. Sutou, T. Omori, K. Yamauchi, N. Ono, R. Kainuma, K. Ishida, *Acta Mater.* 53 (2005) 4121-  
17 4133.
- 18 [53] S. Alkan, H. Sehitoglu, *Acta Mater.* 175 (2019) 182-195.
- 19 [54] P. Sedmak, J. Pilch, L. Heller, J. Kopecek, J. Wright, P. Sedlak, M. Frost, P. Sittner, *Science* 353  
20 (2016) 559-562.
- 21 [55] J.A. Shaw, S. Kyriakides, *Int. J. Plast.* 13 (1998) 837-871.
- 22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65

**Declaration of interests**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

The authors declare the following financial interests/personal relationships which may be considered as potential competing interests:



## Cover letter

Dear Editor,

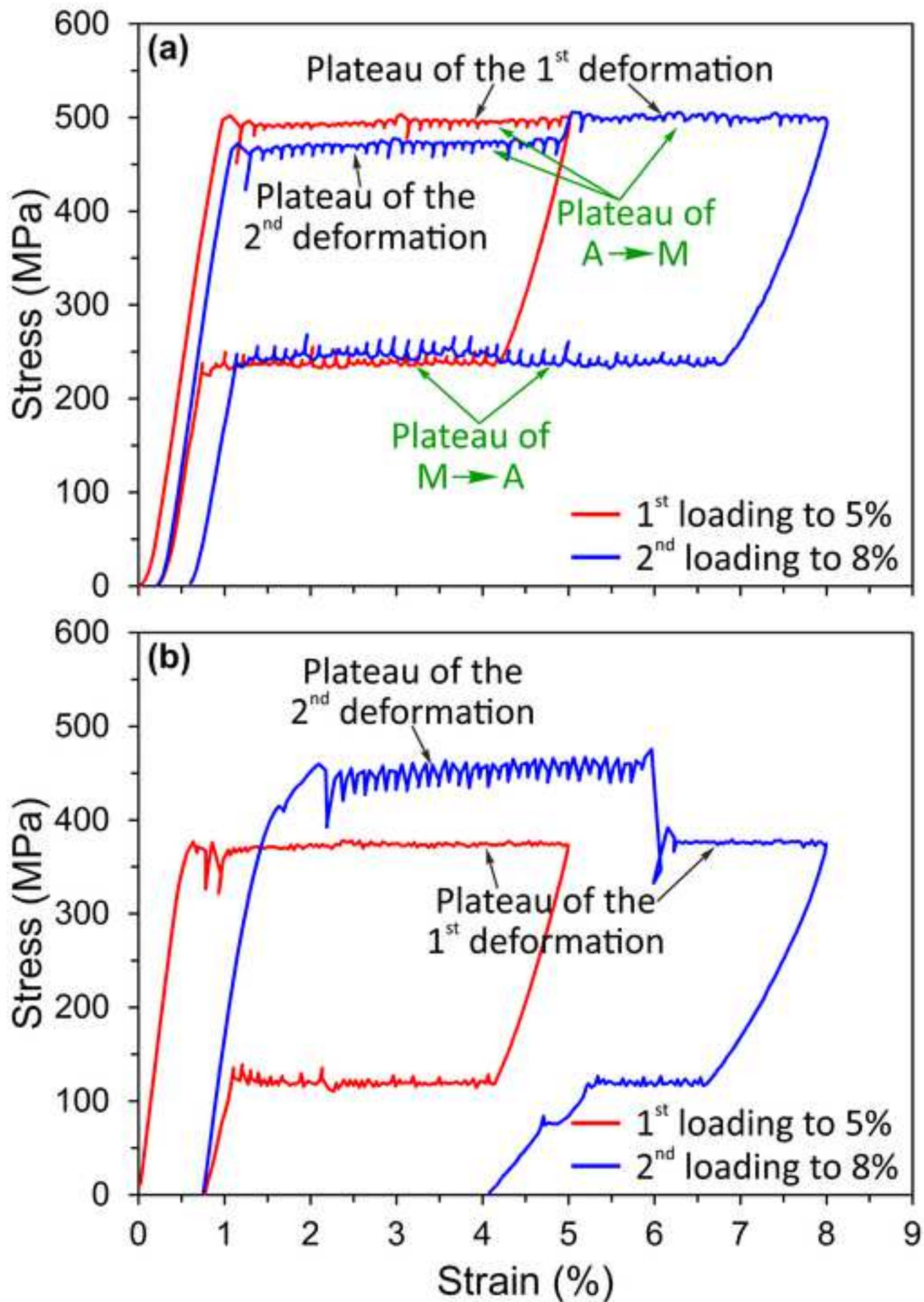
Thank you very much for your editorial work with our submission (**SMM-21-0686**) entitled **“Anomalous stress-strain behavior of NiTi shape memory alloy close to the border of superelastic window”**. This manuscript is the revised version of the submission.

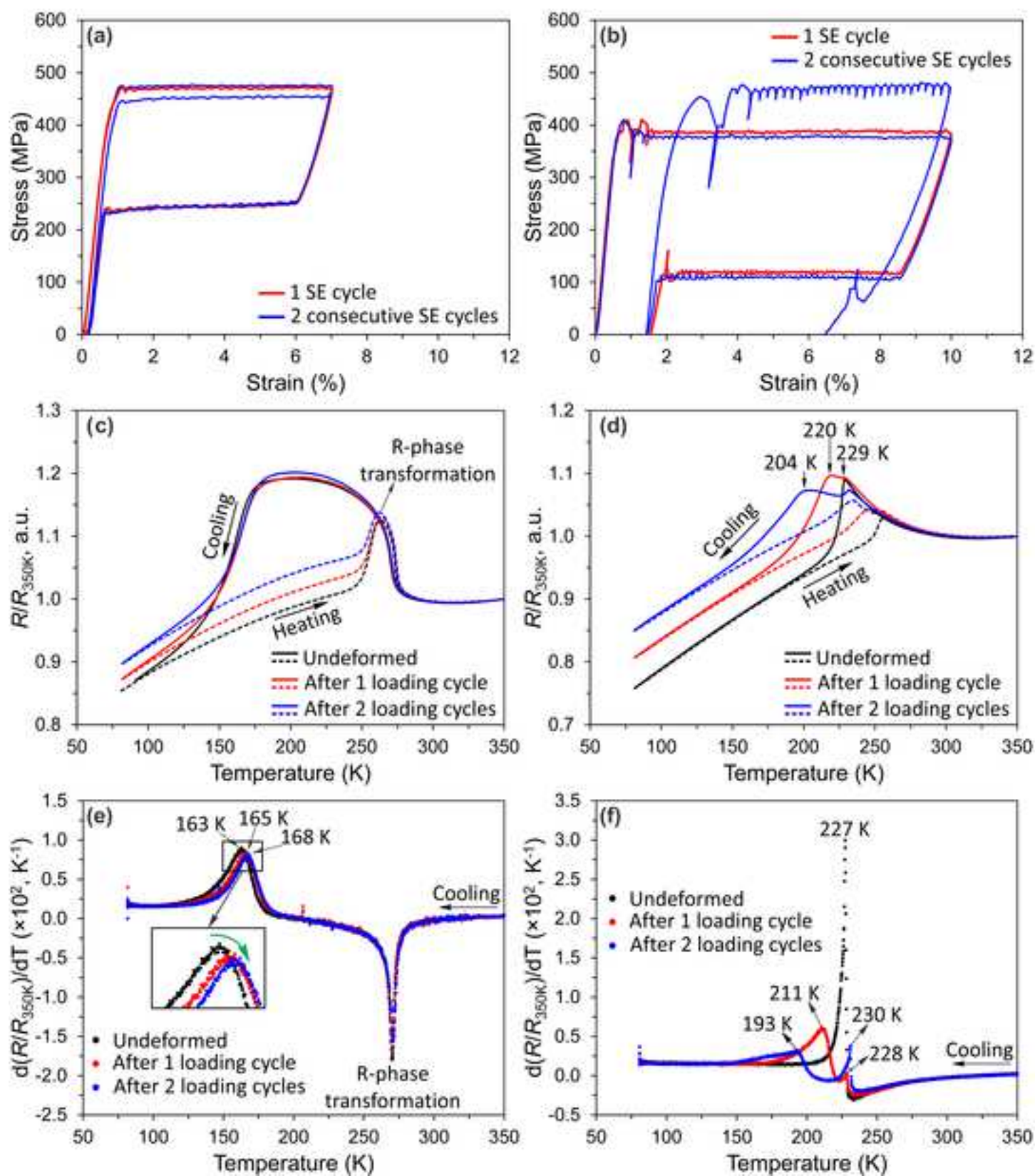
Based on the reviewers' suggestions, we made some changes in our revised manuscript. The changes are highlighted with yellow background (**black text on yellow**). The detailed response to the reviewers' suggestions is shown in “Response to reviewers' comments”.

Yours sincerely,

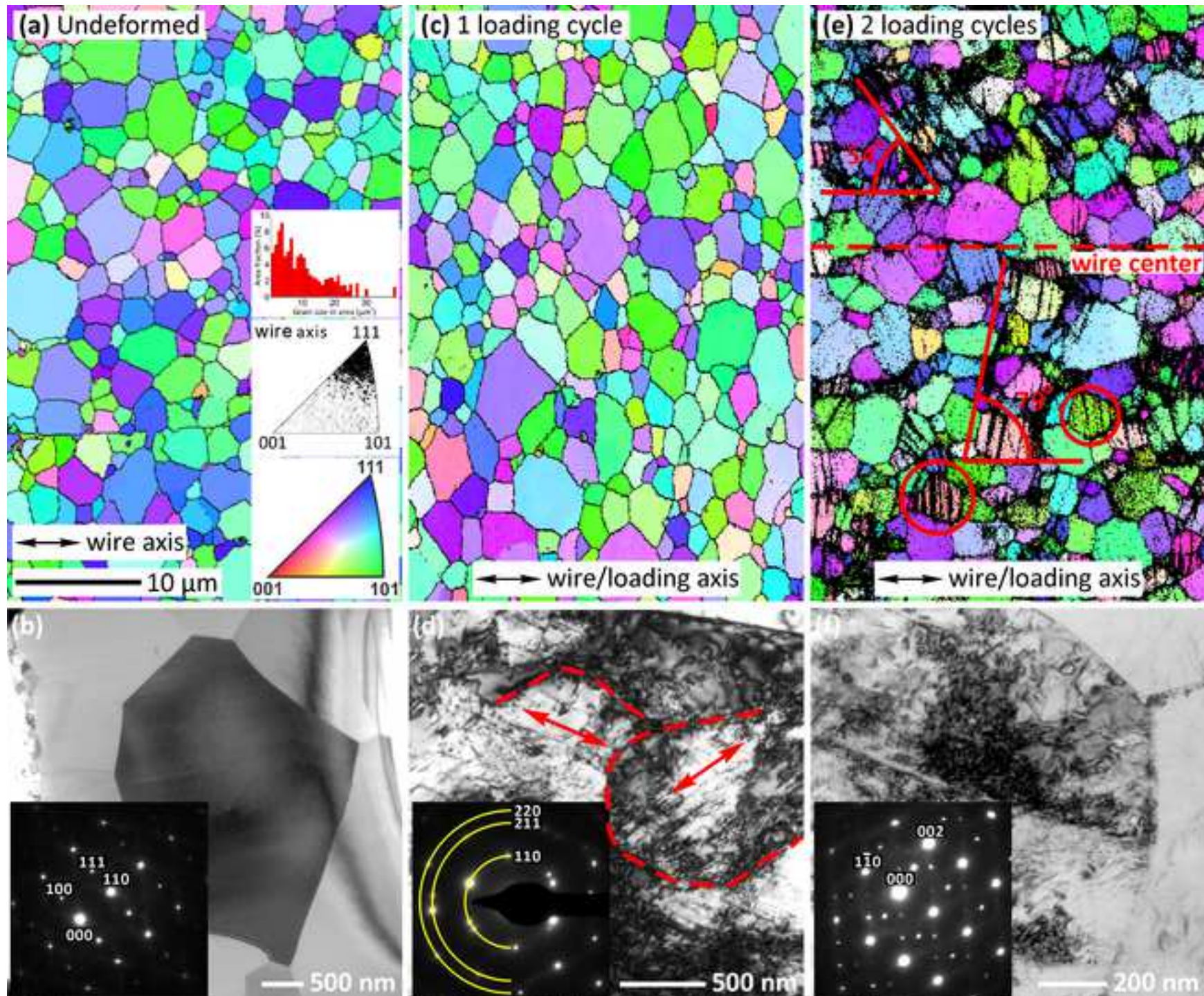
Xiebin Wang

(On behalf of all co-authors)

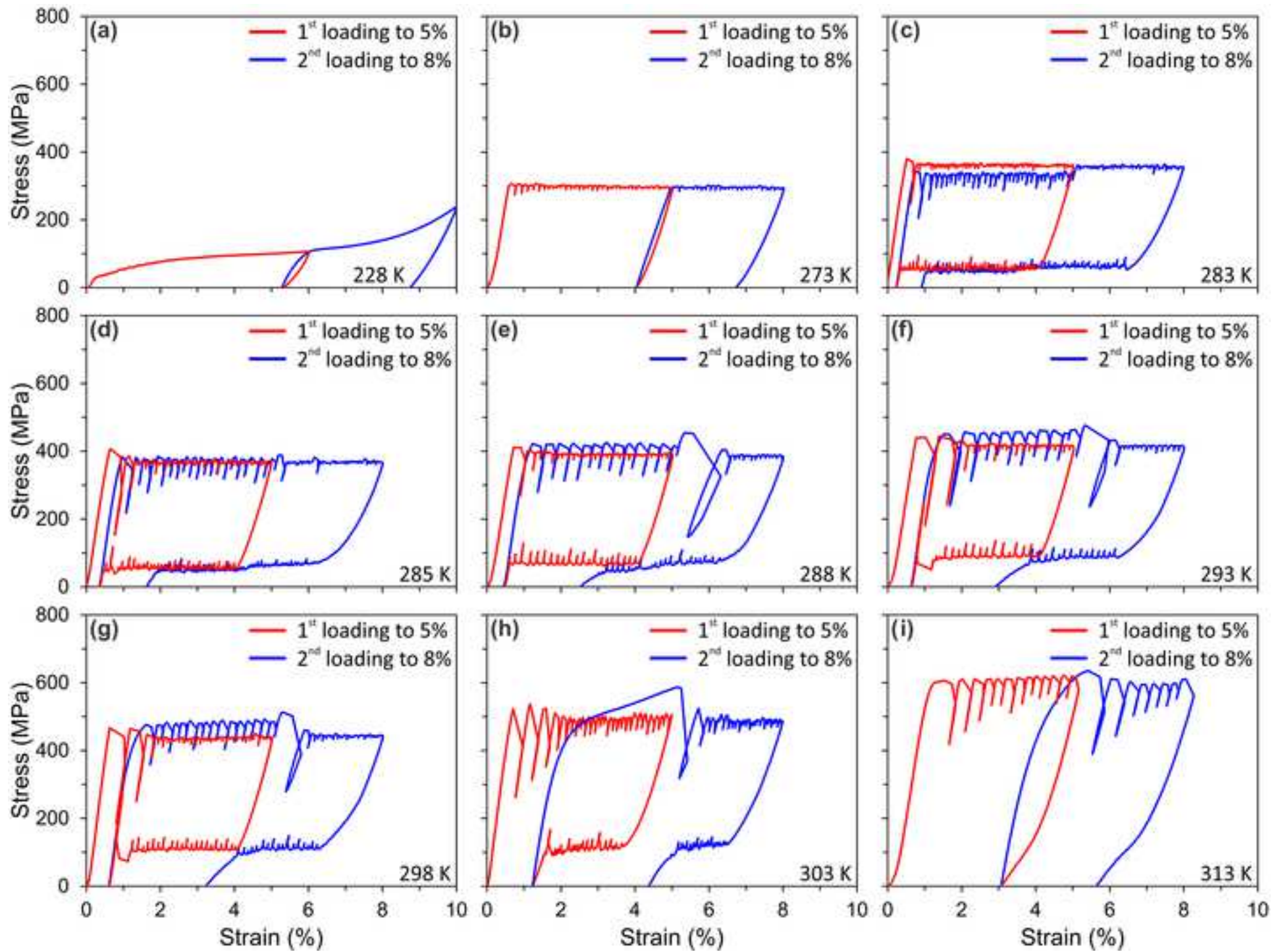


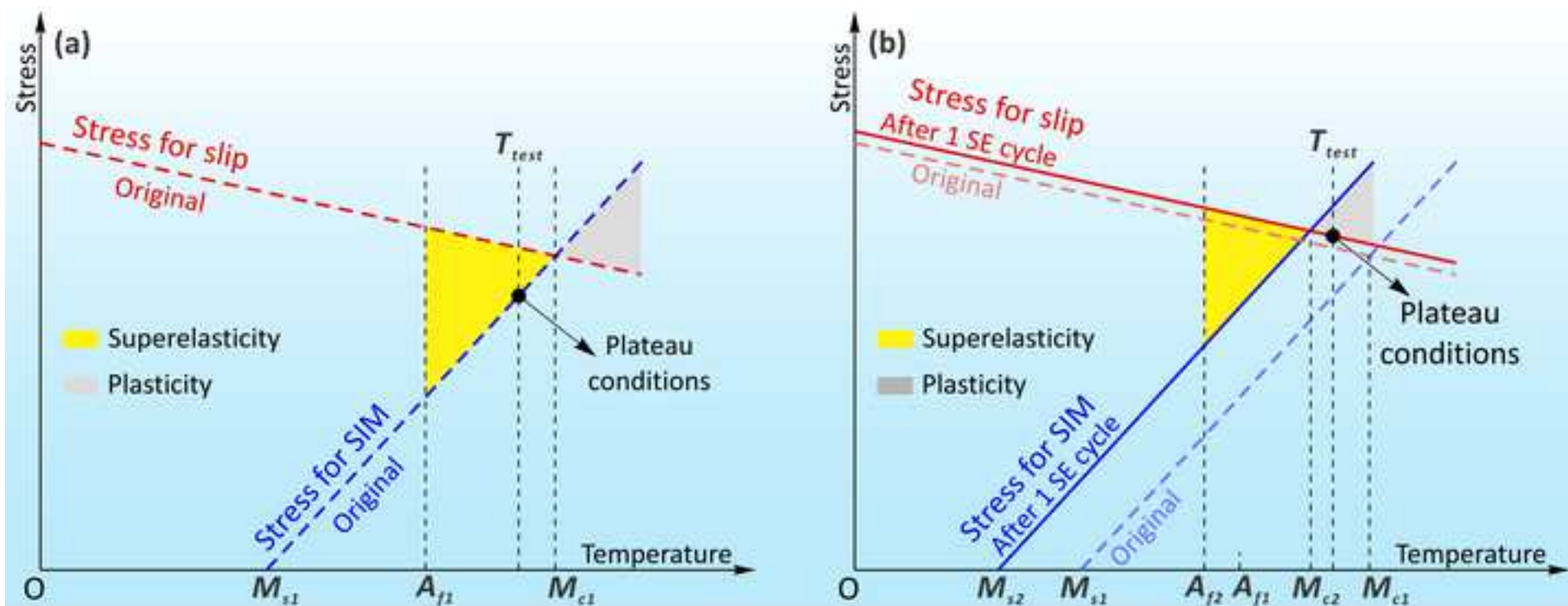














Click here to access/download  
**Supplementary Material**  
Supplyment.docx

