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# Producing oxygen and fertiliser with the Martian atmosphere using microwave plasma

Seán Kelly<sup>a\*</sup>, Claudia Verheyen<sup>a,b</sup>, Aidan Cowley<sup>c</sup> and Annemie Bogaerts<sup>a</sup>

<sup>a</sup>Research group PLASMANT, Department of Chemistry, University of Antwerp, Universiteitsplein 1, BE-2610 Antwerp, Belgium <sup>b</sup>Chimie des Interactions Plasma-Surface (ChIPS), CIRMAP, Université de Mons, 23 Place du Parc, 7000 Mons, Belgium <sup>c</sup>European Astronaut Centre, European Space Agency, Linder Höhe, D-51147, Germany

#### **SUMMARY**

The potential of microwave (MW) plasma-based in-situ utilisation of the Martian atmosphere is explored with focus on the novel possibility of fixing  $N_2$ , for fertiliser production. Conversion in a simulant plasma (i.e., 96/2/2 %  $CO_2/N_2/Ar$ ), performed under energy conditions similar to the Mars oxygen in-situ experiment (MOXIE) currently on-board NASA's Perseverance rover, demonstrates that  $O/O_2$  formed through  $CO_2$  dissociation facilitates the fixation of the  $N_2$  fraction via oxidation to  $NO_x$ . Promising production rates for  $O_2$ , CO and  $NO_x$  of 47.0, 76.1 and 1.25 g/h, respectively, are recorded with corresponding energy costs of 0.021, 0.013 and 0.79 kWh/g. Notably,  $O_2$  production rates are ~30 times higher than those demonstrated by MOXIE, while the  $NO_x$  production rate represents a ~7% fixation of the  $N_2$  fraction present in the Martian atmosphere. MW plasma-based conversion therefore shows great potential as an in-situ resource utilisation (ISRU) technology on Mars, simultaneously fixing  $N_2$  and producing  $O_2$ .

Keywords: plasma; nitrogen fixation on Mars; Mars; In-situ Resource Utilisation; ISRU Mars; Oxygen on Mars; plasma-based gas conversion.

#### INTRODUCTION

The expansion of in-situ resource utilisation (ISRU) technology  $^{1-8}$  will be a key enabler for both private and public funded space exploration of planets like Mars. The prohibitive cost of bringing fuel,  $O_2$  and food to Mars greatly motivates exploitation of local resources. ESA's Ariane 5G heavy lift rocket, for instance, has a payload cost of  $\sim 10,000 \text{ s/kg}$  to reach low Earth Orbit (LEO)  $^9$ . The additional costs of sending each kg from LEO to Mars is estimated at over ten times the initial LEO costs  $^{10}$ , leading to a conservative estimate of 100,000 s/kg. This is further compounded by the  $\sim 26$ -month launch window between the Earth and Mars (i.e., the Hohmann transfer orbit), requiring significantly more resources for both human and two-way robotic missions to the red planet compared to what was needed during the Apollo Lunar missions (taking just over eight days from lift off to splash down). Such staggering costs emphasise the necessity for gaining resources in-situ by bringing the means of production rather than traditional supply-dominated payloads. Leveraging technologies for ISRU by 'living off the land' is therefore a central tenet for future space enterprises, with hopes to efficiently utilise locally available renewable electricity, such as solar, to harvest

<sup>\*</sup>E-mail address: sean.kelly@uantwerpen.be, lead contact

and process native resources. The ISRU paradigm therefore opens many new possibilities for future space exploration <sup>11,12</sup>.

In 2021 the National Aeronautics and Space Administration's (NASA) Martian rover 'Perseverance' performed the 'Mars  $O_2$  in-situ resource utilisation experiment' (MOXIE)  $^{13}$ , producing for the first time, extraterrestrial  $O_2$  using solar-harvested electricity. This milestone event is set to expand innovation in technologies to harvest Martian resources for fuel, life support and materials over the coming decades. The solid oxide electrolysis cell (SOXE) component of the MOXIE experiment demonstrated the production of  $\sim$  6g of  $O_2$  from compressed Martian ambient  $^{14,15}$  using a full sol energy allocation of 1 kWh (further details are given in the **supplemental information** (**figure S1**)). Human consumption of  $O_2$  is about 1 kg/day  $^{16,17}$ , while utilisation of  $O_2$  in a fuel mixture (e.g., Methalox) to power a Mars Ascent Vehicle (MAV) could require thousands of tonnes of fuel  $^{14,18,19}$ . Clearly, significant scale-up is required to achieve the aspirations of future ISRU based missions. MOXIE is based on solid oxide electrolysis cell technology and due to the long start-up requirements ( $\sim$ 2 hours), the technique is largely inflexible to fluctuating energy production and so requires battery storage from any local renewable energy harvested on Mars. Techniques such as plasma-based gas conversion  $^{20,21}$ , which can match production with the availability of renewable electricity (i.e., fast start-up time), would therefore hold great potential for ISRU applications on the red planet.

Mars, with an atmosphere of ~96% CO<sub>2</sub>, ~2 %  $N_2$  and ~2 % Ar  $^{22,23}$ , provides quite favourable low pressure (~1 % of the Earth atmosphere) and temperature (~o °C to < -60 °C) conditions <sup>22,23</sup> for efficient plasma conversion compared to Earth bound climes. To date, however, the potential for this enticing technology remains largely unexplored in an ISRU context <sup>24-31</sup>. Plasma-based O<sub>2</sub> generation and membrane extraction has previously been proposed using direct current (DC) plasma by Wu et al. 24. Gruenwald et al. 25,26 envisaged the use of plasma technology by early Martian settlers for a wide range of applications including O2 production. More recent reports by *Guerra et al.* <sup>27,28</sup> have demonstrated, again, the potential of DC plasmas under Martian conditions, reiterating the benefits of using the ambient conditions for O<sub>2</sub> production. Premathilake et al. reported on the use of a DC plasma generated in-situ of a thin silver membrane which enabled partial oxygen removal <sup>29</sup>. Moses et al. <sup>30</sup> uniquely suggested the harvesting of plasma-produced O<sub>2</sub> during the landing descendent to Mars, where solid oxide cell technology could be incorporated into the heat shield to capture O<sub>2</sub>. The potential use of microwave (MW) plasma under Martian conditions was recently explored in the PEOMA ('Plasma Extraction of O<sub>2</sub> from Mars Atmosphere') project by Wheeler et al., supported by NASA 31. Their experimental study, which focused on O2 production, showed the feasibility for high levels of CO<sub>2</sub> conversion under low Martian pressure (i.e., with low gas flow rate). The energy efficiencies reported via conference proceedings <sup>31</sup> were, however, quite low (i.e., < 10 %). This contradicts previous studies of CO<sub>2</sub> plasma conversion at low pressures (i.e., close or below Martian ambient conditions), which claimed both high efficiency and high conversion levels, even up to 90% at supersonic flow conditions, as reported in the 1980's 32 although the latter findings have not been reproduced to date 33-35. The proposal for N<sub>2</sub> fixation (NF) with the Martian atmosphere using plasma, has, to the best of our knowledge, not been previously interrogated. This intriguing prospect will therefore be of particular focus in this work.

Artificial  $N_2$  fixation is a cornerstone of modern civilisation and currently sustains over 40 % of Earth's population  $^{36,37}$ . N is a rate-limiting nutrient for plant growth and is notably absent from the Martian regolith, unlike the other key macronutrients for plant growth, such as K and P, which have been discovered by soil sampling  $^{38-43}$ . Recent botanical experiments using Martian regolith simulants  $^{38,41}$  have highlighted that seed germination and plant growth may be possible in controlled environments, such as an underground greenhouse. Future utilisation of Martian regolith as a farming substrate will therefore require production of

 $N_2$ -based fertiliser as a key enabler for plant growth to sustain future habitats  $^{7,38,41,44}$ . Another important potential use for NF on Mars is the production of explosives with potential use for excavations and active seismology studies  $^{45,46}$ .

On Earth, industrial-scale NF is at present achieved via the Haber-Bosch (H-B) process (producing ammonia) in combination with the Ostwald process (converting ammonia to nitric acid)  $^{47}$ . This energy-intensive process, fueled by natural gas (i.e., methane), has dominated artificial fertiliser production for almost a century and has enabled crop yield enhancements, which at present nourish a large proportion of the world population  $^{48}$ . Owing to the exceptional stability of the  $N_2$  triple bond, the H-B process is an energy-intensive chemical process, which accounts for  $\sim$ 2 % of the world's energy consumption, consumes  $\sim$ 3 % of the global natural gas output and as a result emits more than 300 million tonnes of  $CO_2$  annually  $^{48}$ . Recently, efforts to find alternatives to the H-B process, which are not reliant on fossil fuels, have expanded, including significant interest in plasma technology  $^{20,36,37,49}$ .

The large-scale use of plasmas to fix  $N_2$  on Earth for fertiliser production goes back to the Birkeland-Eyde (B-E) process to produce nitric acid and was first developed in the early twentieth century  $^{50-52}$ . The B-E process converted air to  $NO_x$  in an electric arc formed inside an electromagnet followed by an oxidation stage where the remaining NO was converted to  $NO_2$  in settling tanks (i.e., a relatively slow reaction preferable at reduced temperatures). This stage was followed by  $NO_2$  hydrolysis in large water absorption towers packed with quartz segments, eventually producing a solution of  $HNO_3$ . The industrialised B-E process produced  $\sim 2\%$   $NO_x$  with an energy consumption for the  $NO_x$  plasma synthesis stage of 2.4-3.1 MJ/mol. The absorption stages added approximately 30-40% additional energy overhead  $^{47,48,51}$ .

A revival of research interest in plasma-based  $N_2$  fixation (i.e., a 21<sup>st</sup> century B-E process) has occurred recently with the expanding availability of renewable electricity and intense efforts to mitigate anthropogenic climate change  $^{48,53}$ . Recently  $^{36}$ , we showed very promising metrics for plasma  $NO_x$  production and energy cost (under Earth conditions), reaching 3.8 % total  $NO_x$  concentration, at a production rate of 0.77 L/min for  $\sim$ 2 MJ/mol energy cost, using atmospheric pressure MW plasmas. To our knowledge, this is the lowest energy cost reported in literature for atmospheric pressure plasmas, at significant  $NO_x$  concentrations. Note that we incorporated modern advancements in MW technology, employing a solid-state amplifier, to power the reactor. Of course, such contemporary technology (and indeed MW technology) was not available during the B-E process development in the early twentieth century  $^{54}$ . Further, solid-state technology can significantly reduce the size and mass associated with MW-powered plasma generation and is therefore much more suitable to meeting the constraints of space deployment  $^{55}$ . Solid-state power supplies are not only more compact but due to superior control (i.e., over frequency and power) enable compact plasma reactors with the need for 'bulky' waveguide components reduced (e.g., plasma lighting applications  $^{56}$ ). A magnetron-based plasma reactor with a volume of  $^{\sim}25$  L and mass  $^{\sim}50$  kg (say) could be reduced to  $^{5}$  L and  $^{5}$  kg  $^{57}$ .

Microwave (MW) generated plasmas offer desirable characteristics  $^{58}$  of high ionisation fraction (i.e., electron density) coupled with relatively low mean electron energies. The electric field generated by the applied MW power selectively heats the electrons, due to their small mass. Furthermore, the electron energy of 1-3 eV, combined with high electron density, gives rise to efficient vibrational and electronic excitation, which can in turn significantly promote efficient (i.e., non-thermal) dissociation routes in gases such as  $N_2$  <sup>53</sup>. Along with  $CO_2$  conversion to  $O_2$  and CO, the potential of plasma-based NF by oxidation of the  $N_2$  content (2 %) available in the Martian ambient provides a potential avenue for NF, which has to date not been explored. Vibrational

and electronic excitation of the  $N_2$  fraction in the Martian atmosphere can lower the threshold for breaking the extremely stable  $N_2$  triple bond (~9.8 eV), which can then be oxidised by  $O/O_2$  formed upon conversion of the large amounts of  $CO_2$  present in the Martian atmosphere.

MW plasma technology offers a very high overall power transmission efficiency compared to low-frequency plasma reactor designs 21. In low-frequency plasma reactors, much of the electrical power can be lost in resistors and reflections due to impedance mismatching (dissipating as heat or nuisance radiation). Even DCpowered plasmas typically have an alternating voltage and current response (due to the RC characteristics of the dynamic plasma and its circuit interaction). Previous studies on plasma-based gas conversion only considered the absorbed plasma power 21 when calculating the energy cost of conversion, which disregards a very large fraction of the real power wasted through transmission to the plasma (e.g., losses can be > 50 % of the overall system power for poorly matched designs, using high resistance components to ballast plasma instabilities), but for practical applications, the input power should be considered, thus accounting also for the transmission efficiency of the power supply. MW plasma using solid state technology can now sustain transmission efficiencies of > 70 % with x10 longer lifetimes compared to magnetron technologies, which degrade in performance comparatively quickly over time. The electromagnetic shielding of high-frequency MW discharges is also much more straightforward compared to low-frequency plasma devices, which present a considerable challenge in terms of minimising electromagnetic interference (EMI) to neighboring electronics (e.g., EMI from arc welders operating at DC and kHz AC are a common nuisance 59,60). We, therefore, believe that MW plasma generated using state-of-the-art low mass, small footprint solid-state amplifiers is the 'technology of choice' for future missions deploying plasma-based gas conversion technologies to Mars and beyond.

In this work, we explore the feasibility of NF on Mars by in-situ leveraging the indigenous atmosphere in combination with the co-conversion of  $CO_2$  to CO and  $O_2$ , for generating resources for fertilisers, fuels and life support systems by means of experiments, benchmarked against MOXIE operating conditions, and supported by chemical kinetics modelling to reveal the underlying mechanism.

### **RESULTS AND DISCUSSION**

The aim of this investigation is to highlight the potential of MW plasma technology for ISRU on Mars. In particular, besides  $CO_2$  conversion into CO and  $O_2$ , we also show the novel possibility for  $N_2$  fixation, the most energy-intensive aspect of fertiliser production and therefore, a key requirement for nourishing any potential future Martian settlers. In the first section, we present the results for our MW plasma experiments using a Martian atmosphere mixture, while in the second section we use our corresponding numerical modelling to reveal the underlying mechanisms. This is followed by a discussion outlining possible utilisation scenarios for plasma-based ISRU on Mars. Finally, a conclusion is given.

# MW plasma-based conversion in a Martian atmosphere

The operating conditions for our MW plasma reactor are inspired by the current MOXIE experiment on board NASA's Perseverance rover <sup>14</sup>. MOXIE operates with an energy allocation of 1 kWh using a full sol worth of solar harvested electricity. The solid oxide cell apparatus (known as SOXE), which performs the dissociation and purification, specifies an operating pressure ranging from 260 to 760 Torr (or 0.34-1 bar). The lower range of 0.34 bar is incorporated here for our plasma operating conditions. In our setup a vacuum system is employed to lower the pressure inside the reactor to ~0.34 bar, while using a typical swirling mass flow rate

of 10 L/min required to operate the reactor <sup>36</sup>. MOXIE uses lower mass flow rates, given the operation conditions of the Martian atmosphere, where gases are compressed from the ambient conditions (~0.01 bar) to a higher pressure (i.e., 0.34-1 bar) inside the SOXE compartment. In spite of such differences, the key parameters of pressure and power are comparable in both cases and thus our investigation serves to gain insight into conversion rates and energy cost possible using plasma-based gas conversion technology under rover energy conditions.

In the **supplemental information figure S1**, we show graphical results of the historic MOXIE experiment for O<sub>2</sub> production completed during April 2021 <sup>15</sup>. MOXIE operating at a power of 300 W <sup>14</sup> over 3.3 hrs of operation (i.e., ~12000 s) produced 5.4 g of O<sub>2</sub>. This encompassed a two-hour warm-up period, followed by ~1 hour of O<sub>2</sub> generation. This yields an O<sub>2</sub> production rate of ~1.6 g/h at an energy cost of 0.19 kWh/g. The overall O<sub>2</sub> produced (i.e., 5.4 g) during the test on April 20<sup>th</sup> (Sol 70 of the Martian year) used the 1 kWh energy allocation available to MOXIE from solar electricity harvested by panels on the Perseverance rover. Note that MOXIE operates at 300 W, while our MW plasma operated at 1 kW. However, we compare the two processes in terms of energy usage rather than power. Indeed, the MOXIE results from April 2021 (used here as a benchmark) consumed ~1 kWh of solar energy which was applied at a rate of 300 W for ~12000 s (~3.33 h). In comparison, our plasma operates at an energy deposition rate of 1 kW which should run for 1 hr to use the same 1 kWh energy as MOXIE. Since our comparison of production rates is made based on the energy required per gram of CO, O<sub>2</sub> and NO<sub>x</sub> produced (i.e., kWh/g) the disparate operating powers are not of particular concern.

In figure 1 A we present the production rates of CO,  $O_2$  and  $NO_x$  (i.e., sum of NO and  $NO_2$ ) measured in the plasma exhaust for an inlet mass flow rate of 10 L/min (i.e., 1.131 kg/h  $CO_2$ , 30 g/h  $N_2$  and 21.4 g/h Ar mass flow rate for a Martian mixture in a ratio 96/2/2%), at (absorbed) plasma power of 1 kW. We measured a production rate for  $O_2$  of 47.0  $\pm$  3.9 g/h, while the production rate for CO was found as 76.1  $\pm$  4.7 g/h. The conversion of  $CO_2$  was measured to be 9.4  $\pm$  0.4 %, with corresponding yields for  $O_2$  and CO measured as 5.2  $\pm$  0.2 % and 9.7  $\pm$  0.6 %, hence the latter approximately two times the  $O_2$  value, which is in-line with the stoichiometry of  $CO_2$  splitting. Notably, the production rates, conversion and product yields were adjusted for the gas expansion, which was measured as ~5% increase of the inlet mass flow (i.e.,  $\alpha = 1.05$ ), see equation 4 in the Experimental Procedures section.

From **figure 1** B we observe that the corresponding energy cost for production of CO and  $O_2$  is 0.0129  $\pm$  0.0008 and 0.021  $\pm$  0.002 kWh/g, respectively. Hence, the plasma produces  $O_2$  at an energy cost about an order of magnitude smaller than in the recent MOXIE test, which reported 0.19 kWh/g (see **supplemental information figure S1**), and the  $O_2$  production rate (47.0 g/h) is even ~30 times higher than in the MOXIE test (i.e., ~1.6 g/h). Plasma-based ISRU therefore shows much promise when compared to solid oxygen electrolysis cell (SOEC) production under comparable energy conditions. Note however that the energy gains demonstrated here do not account for the considerable cost of gas compression and separation, along with the energy losses of the solid-state MW power supply employed in this study, which has an efficiency of ~50 %  $^{61}$ ). The energy cost of compressing the Martian atmosphere by the MOXIE experiment, carried out using a scroll pump, is approximately one third of the total energy cost (i.e., ~0.06 kWh/g)  $^{14}$ . The energy cost of separation is more difficult to estimate, but using the heat available in the plasma exhaust ( $T_{exhaust} > 1000$  K, see discussion below) will be key. Utilisation of this energy could enable the use of  $O_2$  separation technologies, which operate at high temperature. This includes emerging technologies such as oxygen transport membranes  $^{62}$ , or indeed more mature techniques such as solid oxide electrolysis cells. Hence,

synergies with such oxygen separation technologies, enabling production of a pure  $O_2$  stream, and plasma technology may yield very interesting pathways for future innovations. The findings presented here therefore strive to highlight and inspire further interrogation of this potential.

Figure 1 A also shows that the total NOx production rate was found to be ~1.25 g/h, with an equivalent concentration in the exhaust mixture of 1320  $\pm$  105 ppm (or 0.13 %). NO and NO $_2$  are the primary species produced with a ratio  $NO/NO_2$  of ~3.4. Notably, no other  $N_xO_y$  species (e.g.,  $N_2O$ ) were detected in measurable quantities. This produced NO<sub>x</sub> concentration corresponds to ~7% of the N₂ inflow being fixed. The corresponding energy cost for NO $_{\rm x}$  production is ~0.79 kWh/g, see **figure 1** B. This is considerably higher than the energy cost of O<sub>2</sub> production, which is like expected given the low fraction of N<sub>2</sub> available in the Martian atmosphere, but at least it serves to give insight into what is possible using plasma-based N₂ fixation on Mars using plasma technology. On Earth, the Haber-Bosch (HB) process fixes about 171 teragram (Tg) of  $N_2$  per year, of which 70 % is used for fertilisers (i.e., 84.7 Tg)  $^{63}$ . The utilisation of this fixed  $N_2$  (albeit very inefficiently <sup>64,65</sup>) supports at least ~2.8 billion people (i.e., ~40 % of the current world population of 7.2 billion) for their food production. The needs of a Martian settler estimated on this (Earth) basis would be ~3 kg N/yr. or 8.2 g/day, and based on the energy cost of 0.79 kWh/g for NO<sub>x</sub> production shown in figure 1, this amounts to ~6.5 kWh/day or ~2.36 MWh/yr. of energy expenditure. The cost of producing the ~1 kg of O₂ required daily per person <sup>16</sup>, needed to sustain life support on Mars, can be similarly estimated as ~20 kWh/day or ~7.3 MWh/yr. using the energy cost of  $\sim$ 0.02 kWh/g for  $O_2$ , presented in **figure 1**. The average daily energy needs for an Earth citizen today are about ~58 kWh 66 with at least 1 % of global energy used for production of nitrogen-based fertilizer (i.e., ~0.58 kWh) <sup>67</sup>. Future Martian settlers will have much greater energy needs compared to their Earth neighbours. Based on our estimates, as detailed above (i.e.,  $\sim$ 6.5 kWh/day for NO $_{\rm x}$ (fertilizer) production and ~20 kWh/day for O₂ production (life support)), nitrogen fertilizers and O₂ alone will require ~27 kWh/day, i.e., almost half the total average Earth residents use today (around 58 kWh, see above).

We measured the exhaust temperature from the plasma using a k-type thermocouple positioned at ~10 cm outside the plasma and afterglow region (i.e., inside the gas connector shown in **figure 3 B** on the right-hand side). We recorded a steady state temperature of ~782 °C (or ~1055 K) for 1 kW absorbed MW power. Note, our gas analysis is performed after the plasma reactor is in operation for at least 15 minutes, after which time the exhausted gas temperature has stabilised, indicating the reactor has reached a steady state operation, although production rates of CO,  $O_2$  and  $NO_x$  stabilise much faster (~10 s) compared to temperature. Utilisation of this heat downstream presents an opportunity to reduce overall energy costs, when this energy could be efficiently recovered <sup>68</sup>, for example, to heat the incoming gas before plasma conversion (with benefits to  $CO_2$  and  $N_2$  dissociation) or to provide heat energy for another chemical reaction or system downstream. In addition, we plan to explore synergies with  $O_2$  separation technologies which operate at high temperatures, such as SOEC <sup>8,14,69</sup> or oxygen transport membranes <sup>62</sup>, by using the heated plasma exhaust to activate a downstream product separation.

# Numerical modelling

We applied our chemical kinetics model to our experimental conditions presented in previous section, and **figure 2** A presents the calculated CO and  $O_2$  concentrations as a function of position in (and after) the plasma region. The measured concentrations (downstream, hence after the plasma) are indicated with stars, for comparison. An  $O_2$  concentration of 6.2 % is predicted at ~35 cm, 15 cm outside the plasma region, which extends from 0 to 20 cm. This compares reasonably well to the value of 5.2  $\pm$  0.5 % measured in our experiment (see star in **figure 2**). The CO concentration predicted by our model is approximately twice the

 $O_2$  concentration, with a value of 12.5 % at ~35 cm, as shown in **figure 2**. This value also compares reasonably to our experimental measurement of 9.7  $\pm$  0.9 %.

Because of this reasonable agreement between model and experiments, we can use the model to analyse the dominant CO and O<sub>2</sub> production reactions. Averaged across the simulation domain, both CO and O<sub>2</sub> are primarily formed via reaction 13, involving the collision of O atoms with CO2, yielding CO and O2. Our analysis reveals that 85 % of the CO and O2 produced via reaction 13 involves the symmetric stretch and bending vibrational modes of CO<sub>2</sub>, namely CO<sub>2</sub>(V<sub>a</sub>-V<sub>d</sub>) (ranging 0.08-0.33 eV in energy) <sup>70</sup>, while 8 % occurs from ground state  $CO_2$  and 7% originates from  $CO_2(V_{1-21})$ . This behaviour is generally in agreement with previous reports on CO<sub>2</sub> conversion under similar (warm plasma) conditions <sup>58,71-73</sup>. Note that our model of course depends on input data, such as chemical reactions and corresponding rate coefficients and cross sections, but as this dominant CO2 conversion mechanism is in agreement with literature, we believe our assumed chemistry is reliable. Other model assumptions and input data in the model, such as mass flow rate, pressure, reactor dimensions, and power density are matched to our experiments as closely as possible. For instance, the power density is determined by the plasma power divided by the plasma volume, both obtained from the experiment, and radial variations are incorporated to represent gas intersecting different regions of the plasma filament, as in our previous work <sup>36</sup>. The asymmetric stretch mode gives rise to the most efficient CO<sub>2</sub> dissociation pathway through a ladder climbing mechanism. However, we see that this mode is relatively supressed here (i.e., ~7 % contribution) due to the relatively high pressures and temperatures under study (0.34 bar, and 2000-3000 K in the plasma filament), which serve to strongly depopulate CO₂(V) via VT relaxation. Lower pressure conditions would allow a much larger contribution of the asymmetric stretch mode, which due to its vibrationally higher energy can significantly reduce the energy cost for dissociation and thus provide improved efficiency. Pressure and temperature conditions more resemblant to the Martian ambient (i.e., ~o.o1 bar, -6o C) should, therefore, serve to further increase the promising metrics discovered here. Since compressing the gas has an energy cost, it would be interesting to understand the influence of lower pressures on the results. Typically, plasma-based CO<sub>2</sub> conversion and N<sub>2</sub> fixation are more energyefficient at lower pressures, because of more pronounced vibrational-translational non-equilibrium, i.e., the vibrational levels of  $CO_2$  and  $N_2$ , which give rise to the most efficient conversion  $^{20,21}$ , are more overpopulated at lower pressure, because of reduced losses upon collisions with the ground state molecules. Our preliminary calculation results indeed confirm this better performance at lower pressure, but these results could not yet be compared with experiments, and are thus somewhat speculative. In our future work, we will investigate the effect of Martian pressures on the performance of both  $CO_2$  conversion and  $N_2$  fixation.

Figure 2 B shows our model predicts a steady state total NO $_{\rm x}$  concentration of ~311 ppm outside the plasma region, at x ~35 cm. Comparatively, the total NO $_{\rm x}$  concentration in our experiment is measured as ~1320 ppm, hence our model clearly underestimates the NO $_{\rm x}$  production. Ramakers et al. <sup>74</sup> studied a CO $_{\rm z}$ /N $_{\rm z}$  gliding arc plasma under similar conditions, and for the lowest N $_{\rm z}$  fraction investigated (5%), they reported a NO $_{\rm x}$  concentration of about 1500 ppm, hence very similar to our results. In addition, they also showed some discrepancy in absolute values between the simulated and measured NO $_{\rm x}$  concentration, but the trend as a function of N $_{\rm z}$  fraction was correctly captured by the model. We believe the discrepancy in absolute values is attributed to the low N $_{\rm z}$  levels simulated, the complexity of the underlying chemistry (which relies on thousands of empirical reaction rate data) and also the inherent physical assumptions required in quasi-1D models of this kind. In spite of this, we believe we can use the model to gain valuable insight into the underlying mechanisms.

We analysed the dominant NO production reactions, given its importance as the key NO<sub>x</sub> species (i.e.,

measured NO/NO $_2$  ratio of ~3.4, and NO is precursor for NO $_2$ , see **reaction 11**). The Zeldovich reaction between N $_2$  and O atoms (**reaction 9**) is found to contribute for 51 % to the overall NO formation. This involves a contribution of 42.5% from electronically excited N $_2$ (E) (mainly N $_2$ (A $^3$  $\Sigma$ ) at 6.2 eV, and N $_2$ (B $^3$  $\Pi$ ) at 7.4 eV), and a contribution of 8.5 % from the vibrationally excited N $_2$ (v) (mainly from the levels v=10-14 which are near the threshold energy for N $_2$  oxidation). The other Zeldovich reaction, **reaction 10**, is found on average to contribute ~18 % to NO formation. The remainder NO is derived from non-Zeldovich reactions, including the reaction of N and its electronically excited state N( $^2$ D) with CO $_2$ (CO $_2$ (V $_3$ -d) (contributing for ~20 %), while three-body reactions of N and O atoms with CO $_2$ /CO $_2$ (V $_3$ -d) account for ~11 %. In summary, the majority of NO formation is found to occur by a combination of the electronically and vibrationally enhanced Zeldovich reactions (**reaction 9** and **reaction 10**), with a total contribution of nearly 70%. This is beneficial, because the electronically and vibrationally excited levels lower the energy threshold of these reactions, thus contributing to energy-efficient NO $_x$  formation. The Zeldovich reaction scheme for NO $_x$  production, as predicted here, is a well-known reaction pathway for nitrogen oxidation. Indeed, this mechanism is consistent with our earlier works in both O $_2$ /N $_2$  mixtures  $^{36,75-78}$  and CO $_2$ /N $_2$  mixtures  $^{74}$ , giving us confidence in the model's capability to describe the chemical mechanisms, in spite of the low NO $_x$  levels predicted.

Finally, our model also provides information on other key plasma parameters, such as electron density and electron temperature. Averaged over the plasma region (i.e., o-20 cm), they were found to be  $3.3 \times 10^{12}$  cm<sup>-3</sup> and 0.51 eV, respectively. This is consistent with previous experimental reports on MW plasmas under similar conditions  $^{35.79}$ .

#### Utilisation scenarios

Future progress of plasma-based ISRU on Mars will need to capitalise on the promising metrics demonstrated here, but will also need to solve outstanding hurdles, by integrating any plasma conversion process with efficient gas separation technologies for a particular utilisation scenario 2. Adsorption techniques (i.e., pressure or temperature swing adsorption) for both CO and O<sub>2</sub> separation from O<sub>2</sub>/CO/CO<sub>2</sub> and CO/CO<sub>2</sub> mixtures are an important separation technology in a Martian ISRU context 80. Carbajo et al. 81 recently investigated the use of zeolite materials for separation of a typical plasma produced CO/O<sub>2</sub>/CO<sub>2</sub> mixture using pressure swing adsorption (PSA). The authors showed that existing commercial materials should perform well (i.e., give reasonably high purity ~96 %) under mild conditions (i.e., ~2 bar). The current energy cost of compression for MOXIE is about ~1-2 Wh/g CO<sub>2</sub> 14,69,82 (pressure range 0.34-1 bar), so a downstream PSA system to separate CO or O<sub>2</sub> would likely have an energy cost on this order (i.e., ~1-2 Wh/q CO, O<sub>2</sub>). In our case the gas inlet flow rate is ~1100 g/h CO2, so the energy costs for compression for a downstream PSA stage would likely be similar to the plasma conversion costs (i.e., per g O₂ or CO produced), yielding a total energy cost, including separation, in the order of ~0.04 kWh/g for O<sub>2</sub> production, or ~0.026 kWh/g for CO production, based on the promising metrics reported here. Hence, these combined costs could still fall significantly below MOXIE. However, this combination of plasma-based CO₂ splitting and PSA still needs to be tested in practice, to confirm these numbers, which are now only theoretical estimates. Furthermore, these separation methods do not account for the nitrogen fixation part. A notable downside of adsorption techniques like PSA is the difficulty of attaining very high purity (i.e.,  $> 96\% O_2$ ), which may be of concern for  $O_2$  utilisation in life support. Several consecutive adsorption and regeneration cycles would be necessary to increase purity 83, leading to higher energy costs. However, given the estimates here, this could still be quite competitive with MOXIE. Further, the ambient low-pressure conditions on Mars could provide a 'free' pressure differential (i.e., the technique known as vacuum pressure swing adsorption (VPSA)) to further reduce this energy overhead. Indeed, a MW plasma can be sustained across a wide pressure range (i.e., from ambient Martian pressures to

above Earth ambient pressures) and deposits excess energy (not used in chemistry) as heat, providing a thermal energy source. This operational flexibility could be valuable for incorporation into any multi-stage adsorption configuration, where combinations of temperature swing adsorption (TSA), PSA and VPSA could provide a rapid separation process (e.g., ~10-100 s) <sup>84</sup> amenable to coupling with intermittent (solar) electricity.

Combinations of plasma and solid oxide electrolysis (SOXE), i.e., technology of MOXIE which decomposes/electrolyzes  $CO_2$  and separates out the  $O_2$  product  $^{14,69,82}$ , could also offer an intriguing prospect for production of highly purified O2. Pandiyan et al. 85 recently interrogated the electrolysis of a CO2 MW plasma-exhaust mixture consisting of CO, O<sub>2</sub> and CO<sub>2</sub>. The authors showed that O<sub>2</sub> separation, in particular, can be achieved exclusively at low overpotentials (i.e., < 0.75 V) and reduced temperatures (i.e., ~650 °C); conditions where CO oxidation and CO<sub>2</sub> electrolysis (i.e., dissociation of CO<sub>2</sub>) are not active. The authors reported a promising energy reduction of > 50 % for O<sub>2</sub> production (i.e., separation) using this combined plasma-SOXE approach, compared to a pure CO<sub>2</sub> SOXE O<sub>2</sub> production. Feeding the SOXE cell with a CO/O<sub>2</sub>/CO<sub>2</sub> (plasma) mix rather than pure CO<sub>2</sub> is also found to benefit considerably the cells durability, with similarities to the methodology of partial recycling of exhaust CO used in MOXIE 14. Notably, any plasma-SOXE hybrid technology operating at reduced temperatures (i.e., < 650 °C) would be much more flexible to powering by intermittent electricity sources and better exploit the operational flexibility of plasma conversion. MOXIE is currently limited in the number of thermal cycles, possibly due to material degradation 14,69, so it cannot be easily switched on and off. On-going research efforts into low-temperature SOXE materials 86 could also have a significant impact on any future hybrid designs reaching more flexible operating conditions.

Obtaining a useable form of fixed nitrogen for farming on Mars could largely depend on the availability of water (e.g., extracted from Martian clays or regolith <sup>87</sup>). Water reacts readily with NO<sub>2</sub>, a process exploited in the original BE process where NO is first oxidised to NO<sub>2</sub> in a settling tank before hydrolysing it in a washing column to form a nitric acid solution. This could potentially solve the separation problem of the formed  $NO_x$ from the other gas components (CO, O2 and unconverted CO2 and N2). Contemporary advances suggest that the energy cost and size of any adsorption stage could be significantly reduced, e.g., using modern NO<sub>x</sub> absorbents such as BaO <sup>88,89</sup> in combination with temperature/pressure swing adsorption <sup>47</sup>. Nitric acid could be deployed directly for use in a hydroponic 'soil-less' farm, serving as a direct source of nitrates for plant growth 90. Applying an acid solution directly to soils in an open environment could lead to relatively poor uptake by plants due to the liquid's volatility. Soil-based Martian farms, as on Earth, would therefore likely benefit from solid forms of nitrogen fertiliser. Combinations with urine is one possible pathway to form solid nitrogen fertiliser. Mixing nitric acid with urine reacts readily (i.e., exothermically) to form urea nitrate, a solid crystalline material which can be mixed into soils as a fertiliser. Further, the fermentation of urine can produce ammonia gas, aided by the urease enzyme. Bubbling ammonia through nitric acid will readily (i.e., exothermically) produce the solid ammonium nitrate. Indeed, ammonium nitrate is the most common form of nitrogen fertiliser used today on Earth for soil-based farming. Notably both urea nitrate and ammonium nitrate are also powerful explosives with potential utility for excavation activities by future settlers.

In summary, we emphasise that the promise of producing  $O_2$  and CO, while fixing nitrogen on Mars using plasma technology combined with SOXE and/or adsorption methods is, at present, largely conceptual; however, we believe this approach holds much potential for Mars ISRU and we hope that this paper can inspire future research efforts.

## **CONCLUSIONS**

We demonstrate the novel possibility of fixing  $N_2$  on Mars, besides  $CO_2$  conversion to CO and  $O_2$ , using the local Martian atmosphere. Microwave (MW) plasma conversion of the majority  $CO_2$  fraction (~96 %) in the Martian atmosphere results in O atom formation, which enables oxidation of the small  $N_2$  fraction (~ 2%), and thus results in fixation of ~7 % of the  $N_2$  present in a simulant Martian atmosphere. Our MW plasma investigation shows promising  $O_2$ , CO and  $NO_x$  production rates of 47.0, 76.1 and 1.25 g/h at an energy cost of 0.021, 0.013 and 0.79 kWh/g. Using the current energy allocation of 1 kWh available to the Mars Oxygen In-situ Electrolysis (MOXIE) experiment on NASA's Perseverance rover, our MW plasma produces 47.0 g/h of  $O_2$ , which is almost 30 times higher than the current capabilities of MOXIE (1.6 g/h), and at 10 times lower energy cost (0.021 kWh/g vs 0.19 kWh/g). Plasma-based conversion, therefore, shows great potential as a future Martian ISRU technology.

The technology also has the key benefit of a rapid start-up time and is therefore highly flexible to the intermittent availability of Martian solar electricity (i.e., energy storage could be forgone). However, the energy costs reported only consider the plasma process, and do not yet account for the cost of gas compression and separation, for which the combination with SOEC would be very interesting, especially because the hot plasma exhaust gas could activate a SOEC for downstream product separation. We hope this paper can inspire future research in this direction. Indeed, an efficient gas separation technology downstream, in combination with our plasma technology, could generate the pure chemical streams for utilisation as fertiliser, life support and fuel for future robotic and human exploration of the red planet.

#### **EXPERIMENTAL PROCEDURES**

# Resource availability

# Lead contact

Further information and requests for resources and materials should be directed to and will be fulfilled by the lead contact, Seán Kelly (sean.kelly@uantwerpen.be).

# Materials availability

This study did not generate new unique materials.

## Data and code availability

Data from this study are available from the lead contact upon reasonable request.

#### Experimental setup

A schematic of our setup and an image of our MW plasma reactor in operation with a Martian simulant mixture consisting of 96% CO<sub>2</sub>, 2% N<sub>2</sub> and 2% Ar are shown in **figure 3** A and **figure 3** B. The power supply is composed of a collection of laterally diffused metal oxide semiconductor (LDMOS) power amplifiers, from which the output powers are combined in a mixer waveguide (WR340). This waveguide is connected via an isolator and an auto-tuner to a tapered waveguide section, including a 16 mm inner diameter quartz tube. The latter is mounted perpendicularly through a coupling hole (i.e., < 1/4 wavelength in diameter), where the plasma ignition takes place. An auto-tuner, impedance analyser and adjustable short are used to tune the electric field to optimal conditions for electrical breakdown, and to sustain a continuously powered plasma with minimum reflected power (<< 5 %). Tangential gas injection ports coupled with a helical insert allow a

swirl or vortex flow within the quartz discharge tube. Upon ignition, a surface wave sustained mode <sup>91</sup> is generated, with the plasma filament located at the tube centre (see **figure 3 B** and **figure 3 C**). This provides a key benefit by isolating the warm plasma (~2000-3000 K) from the quartz tube walls, allowing for elongated and stable plasma column formation along the tube lateral axis. Once initiated, the surface wave mode is stable across a wide pressure range from low pressures (~0.1 bar) to several bar.

Analysis of the  $NO_x$  species (i.e., NO and  $NO_2$ ) and CO in the exhaust gas is performed using non-dispersive infra-red and ultra-violet (NDIR/UV) absorption spectrometry (Rosemount X-STREAM XEGP Continuous Gas Analyzer  $^{92}$ ), while  $O_2$  is measured using a PyroScience GmbH  $^{93}$  sensor based on an infra-red luminescent quenched absorption technique. All diagnostics were calibrated using pre-mixed calibration gases (Air Liquide) and cross-checked with gas chromatography (GC) using the compact-GC instrument from Interscience. This GC has two channels, each with a Thermal Conductivity Detector (TCD), using carboxen and molsieve columns (1010 PLOT and 5A, respectively) for  $O_2$ ,  $N_2$  and CO detection, and two RT Q-bond columns (3 m and 10 m length, respectively) for  $CO_2$  detection  $^{94}$ .

The primary gas converted in our experiments is  $CO_2$  (96 % of the Martian atmosphere), with oxidation of the small  $N_2$  content (2 %). The argon fraction (2 %) is not converted due to its inertness. The key overall reactions for consideration are therefore:

$$CO_2 \rightarrow CO + 1/2 O_2$$
 (Reaction 1)

$$O_2 + N_2 \rightarrow 2 NO$$
 (Reaction 2)

$$2 O_2 + N_2 \rightarrow 2 NO_2$$
 (Reaction 3)

In any gas conversion process, there is typically gas expansion or contraction due to changes in stoichiometry. In our experiments this results in an increase to the mass outflow. Indeed, given  $CO_2$  is the primary component of the gas fraction, plasma conversion to CO and  $O_2$  results in an expansion of the inlet flow, which depends on the degree of conversion. Notably the formation of  $NO_2$  results in gas contraction, however, due to the relatively small fraction of  $N_2$  and even smaller fraction of  $NO_2$  formed compared to CO and  $CO_2$ , this has negligible effect compared to expansion from  $CO_2$  conversion. Strategies for measuring the gas mass outflow can include direct measurement or inference of the degree of expansion/contraction using dilution gases. In our case, given the dominance of  $CO_2$  conversion, the degree of expansion can be inferred using the  $CO_2$  conversion as:

$$\alpha = 1 + 0.5 * \eta_{\text{CO}_2}^{converted} \tag{4}$$

$$\eta_{\text{CO}_2}^{converted} = \frac{\eta_{\text{CO}_2}^{OFF} - \alpha * \eta_{\text{CO}_2}^{ON}}{\eta_{\text{CO}_2}^{OFF}}$$
 (5)

Where  $\alpha$  represents the gas expansion factor (i.e.,  $\alpha>1$ ),  $\eta_{\rm CO_2}^{converted}$  is the fraction of CO<sub>2</sub> converted,  $\eta_{\rm CO_2}^{OFF}$  is

the fraction of CO<sub>2</sub> in the mixture when the plasma is off (i.e., o.96 in our case) and  $\eta_{\text{CO}_2}^{ON}$  is the measured CO<sub>2</sub> fraction when the plasma is ON. Rearranging **equations 4** and **5** we can solve them to find  $\alpha$  based on the measured CO<sub>2</sub>:

$$\alpha = \frac{1.5 * \eta_{\text{CO}_2}^{OFF}}{\eta_{\text{CO}_2}^{OFF} + 0.5 * \eta_{\text{CO}_2}^{ON}} \tag{6}$$

Subsequently, when knowing  $\alpha$ , we obtain the CO<sub>2</sub> conversion by **equation 5.** The CO, O<sub>2</sub> and NO<sub>x</sub> production rates are calculated based on the percentage yield of each species measured in the exhaust and using the corresponding mass flow rate adjusted for the gas expansion. The individual production rates are then calculated for CO, O<sub>2</sub> and NO<sub>x</sub> (NO +NO<sub>2</sub>) as:

$$PR_{CO,O_2,NO_x}[g/h] = \frac{\eta_{CO,O_2,NO_x} * \mu_{CO,O_2,NO_x}[g/mol] * \alpha * f_{in}[L/min]}{22.4 [L/mol]} * 60 [min/h]$$
(7)

where the production rate, PR [g/h], is determined by the inlet mass flow rate  $f_{in}$  [L/min] (standard litres per minute) for the Martian atmosphere simulant mix, consisting of 96 % CO<sub>2</sub>, 2 % N<sub>2</sub> and 2% Ar.  $\eta_{CO/O_2/NO_x}$  represents the fraction of species produced, directly measured in the plasma exhaust while  $\mu_{CO,O_2,NO_x}$  [g/mol] is the molar mass of CO, O<sub>2</sub> or NO<sub>x</sub>, and 22.4 [L/mol] is the molar volume of a gas under the corresponding standard conditions (i.e., standard temperature and pressure) for which our mass flow controllers are calibrated.

The energy cost for the production of CO,  $O_2$  or  $NO_X$  is then obtained using:

$$EC_{CO,O_2,NO_x}[kWh/g] = \frac{Power[kW]}{PR_{CO,O_2,NO_x}[g/h]}$$
(8)

where the Power [kW or kJ/s] is the absorbed MW power measured during steady state plasma operation.

#### Numerical modelling and chemistry

A quasi-1D chemical kinetics model is employed using the ZDPlasKin (Zero-Dimensional Plasma Kinetics) solver <sup>95-97</sup>. An overview of the simulation scheme is given in **figure 4** with full details of the equations employed found in our earlier work <sup>36</sup>. The time-evolution of the species densities, including electrons, and various charged and neutral species, is calculated by balance equations, considering the production and loss terms by chemical reactions. Dynamic changes in the gas velocity due to temperature and stoichiometric changes in the gas mixture are updated on each time step.

The power density  $P[W/cm^3]$  is derived from our experimental measurements of the absorbed power (i.e., forward minus reflected power) combined with determination of the plasma volume via camera imaging inside the tapered section of the waveguide (see **figure 3 B** and **figure 3 C**). A cylindrical shape is assumed for the plasma volume  $^{98}$ . This is consistent with vortex-stabilised discharges, where the plasma is contained within the tube inner region, separated from the containment walls by a swirling or vortex flow boundary. The plasma elongates along the direction of the flow (i.e., along the axial extent of the reactor tube) forming a cylindrical shape in its steady state  $^{34,98,99}$ . In order to account for the radial variation in power density from

the centre of the plasma filament to its edge, the light emission across the radial extent of the plasma filament at its ignition point inside the waveguide (as shown in **figure 3 C**) is used as a proxy for the plasma width. We therefore solve the quasi-1D model for two different radial sections, with a corresponding high power density of the plasma core and a relatively low power density to represent the plasma edge, as explained in our earlier paper <sup>36</sup>. For each of both quasi-1D models, we assume a triangular distribution of power density in the lateral extent of the plasma (i.e., along the direction of gas flow), in line with earlier modelling of power dissipation in surface wave sustained MW plasmas <sup>58,100</sup>.

The reduced electric field, i.e., the ratio of electric field over gas number density, a key fundamental variable defining the plasma characteristics, is calculated from our measured specified power density (see **figure 4**). A Boltzmann solver (i.e., BOLSIG+) is utilised to simulate electron dynamics by linking the plasma conductivity (a function of the reduced electric field) to the electron mobility. Further to this, the gas temperature is solved in the model at each time step, based on gas heating due to elastic collisions of electrons with the gas molecules, the enthalpy contributions from the chemical reactions between all plasma species, heat losses to the walls and the dynamic heat capacity taking account of the gas mixture. The radially averaged gas temperature (which is assumed to have a parabolic profile) is calculated by considering the time-dependent gas thermal balance equation under isobaric conditions. Further details can be found in our earlier works <sup>36,71,101,102</sup>.

The Martian air chemistry (i.e.,  $CO_2/N_2/Ar$ ) employed here is assembled from our earlier works, where details can be found  $^{71,74,103}$ . 149 species are included in the model, i.e., the electrons,  $CO_2$ , Ar,  $N_2$ ,  $O_2$ , CO, various  $N_xO_y$  molecules in ground state, along with various (vibrational and electronic) excited levels, various radicals, atoms and ions (see **Table 1**). These species react with each other in 973 electron impact reactions and 12,604 heavy particle reactions (i.e., between molecules in ground state or excited levels, radicals, atoms or ions). For the heavy particle reactions, the rate coefficients are adopted from our earlier works, whereas the rate coefficients for the electron impact reactions are calculated using the Boltzmann solver BOLSIG+  $^{97}$  built in 7DPlasKin.

Plasma-based conversion of relatively inert molecules, such as  $N_2$  and  $CO_2$ , provides unique reaction pathways not available in purely thermal conversion. Especially in MW plasma, the electrons have the right energy (~ 1 eV) to cause excitation towards the lowest vibrational levels in  $CO_2/N_2/O_2$ , followed by further vibrational-vibrational (V-V) collisions, which enable a "ladder-climbing" process, gradually populating higher vibrational levels (denoted as e.g.,  $CO_2(V)$  and  $N_2(V)$ ). Further significant populations of electronically excited species (e.g.,  $CO_2(E)$  and  $N_2(E)$ ) can form inside the plasma region. Such species serve to lower the overall energy required for  $CO_2$  dissociation (i.e.,  $O_2$  formation) and  $NO_x$  formation, as their higher energy levels help to overcome the activation barriers.

The underlying elementary reactions for  $NO_x$  production (see overall **reactions 2** and **3** above) in a plasma involve the atoms formed upon dissociation of the corresponding molecules, and proceed via the (electronically or vibrationally enhanced) Zeldovich mechanism  $^{53,104,105}$  consisting of the following reactions:

$$O + N_2/N_2(E, V) \leftrightarrow N + NO$$
 (Reaction 9)  
 $N + O_2/O_2(E, V) \leftrightarrow O + NO$  (Reaction 10)

The above reaction pair is typically rate-limited by **reaction 9**, given the energy requirement for overcoming the strong  $N_2$  triple bond. Notably, the mechanism here can be significantly different from the purely thermal Zeldovich mechanism (i.e., involving only ground state  $N_2$  and  $O_2$ ) due to the presence of vibrationally or electronically excited  $N_2$  and  $O_2$  molecules, available in plasmas. The vibrationally or electronically excited  $N_2$  molecules lower the dissociation threshold required to break the  $N_2$  bond (~9.8 eV) through colliding with O atoms (i.e., **reaction 9**). The N atoms formed in **reaction 9** can then further react with both ground state and vibrationally/electronically excited  $O_2$  molecules (**reaction 10**) to produce another NO. **Reaction 10** also produces an additional O atom, which can again react with ground state and vibrationally/electronically excited  $N_2$  molecules (i.e., **reaction 9**), or oxidise NO to produce  $NO_2$  (**reaction 11**):

$$NO + O \rightarrow NO_2$$
 (Reaction 11)

A similar oxidation pathway applies to  $CO_2$ . The overall **reaction 1** above includes the following elementary reactions, involving atomic oxygen:

$$CO_2(E, V) \rightarrow CO + O$$
 (Reaction 12)

$$CO_2(E,V) + O \rightarrow CO + O_2$$
 (Reaction 13)

Direct dissociation processes in **reaction 12**, such as electron impact dissociation, have an energy threshold of 5.5 eV to overcome the (ground state)  $CO_2$  bond energy, but the atomic oxygen produced in **reaction 12** can further react (i.e., **reaction 13**) with  $CO_2$  (and its vibrationally or electronically excited states). This coupling lowers the threshold considerably (e.g., for ground state  $CO_2$  this lowers the overall energy threshold to 2.9 eV <sup>35</sup>).

### SUPPLEMENTAL INFORMATION

Document 'Supp\_info.pdf' is the main supplemental PDF.

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## **AUTHOR CONTRIBUTIONS**

S.K. conducted the experiments and modelling including data analysis; S.K., A.C and A.B. wrote the paper. C.V. conducted modelling and compiled the chemistry.

# **DECLARATION OF INTERESTS**

The authors declare no competing interests.

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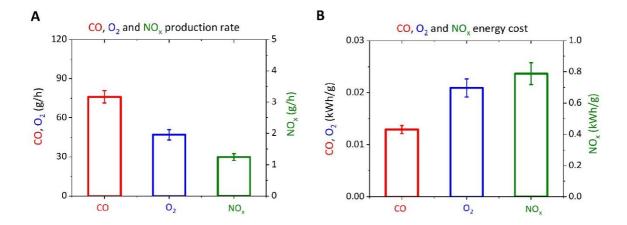


Figure 1: A: Absolute production rates (g/h) of CO,  $O_2$  and  $NO_x$ ; B: equivalent energy cost (kWh/g), in a MW plasma using a Martian simulant mixture of  $CO_2/N_2/Ar$  (96/2/2 %) at 10 L/min flow rate, 0.34 bar pressure and 1 kW absorbed power. Note the  $NO_x$  data is indicated on the right-hand y-axes.

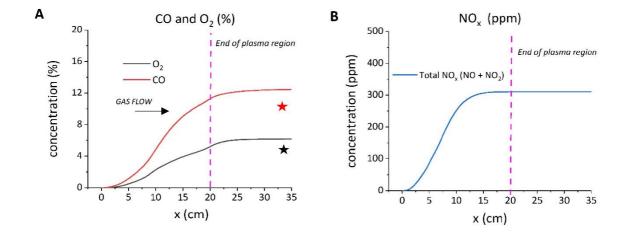


Figure 2: Calculated CO,  $O_2$  (figure 2 A) and  $NO_x$  (figure 2 B) product concentrations as a function of position in and after the MW plasma, using a Martian simulant mixture of  $CO_2/N_2/Ar$  (96/2/2% at 10 L/min flow rate, 0.34 bar pressure and 1 kW absorbed power). The measured CO and  $O_2$  concentrations (downstream, hence after the plasma) are indicated with stars in figure 2 A, for comparison. The measured  $NO_x$  concentration is not added in figure 2 B, because there is still quite large discrepancy with the calculated value (see text). Note: the direction of gas flow through the plasma is along the positive x direction as indicated in the figure.

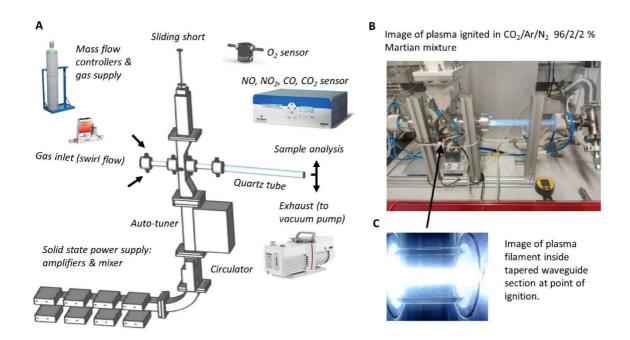


Figure 3: A: Illustration of our MW reactor consisting of a solid-state MW power supply, circulator, autotuner and tapered waveguide section terminated by a sliding short. The plasma is ignited inside a quartz tube where a swirling flow is injected. Sample analysis of the exhaust gas is carried out using an NDIR and a luminescence O₂ sensor. B: Photo of the reactor in operation with a Martian simulant atmosphere at 0.34 bar pressure and 1 kW power. Ignition takes place in a tapered section of a 2.45 GHz WR340 waveguide where a plasma is suspended at the centre of the tube. C: In-waveguide photo of the plasma; camera viewpoint is looking towards the quartz tube inside the waveguide, as indicated by the arrow.

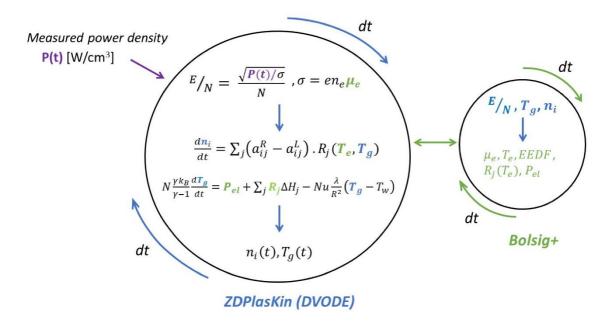


Figure 4: Overview of the numerical solution scheme for the quasi-1D model employing ZDPlasKin, which incorporates the DVODE code for ODE integration and BOLSIG+ to solve the Boltzmann equation at each time step. An experimentally determined power density [W/cm³] is coupled externally to enable solution via the reduced electric field E/N.

Table 1: Species included in our quasi-1D model for a CO<sub>2</sub>/N<sub>2</sub>/Ar mixture.

Neutral	Excited	Charged
$CO_2$ , $CO$ , $C_2O$ , $C$ , $C_2$ , $CN$ , $ONCN$ , $NCO$ , $NCN$ , $C_2N$ , $C_2N_2$	$CO_2(V_{a_1}V_{b_1}V_{c_1}V_{d})$ , $CO_2(V_{1^{-1}}V_{21})$ , $CO_2(E_1)$ , $CO(V_{1^{-1}}V_{10})$ , $CO(E_1-E_4)$	CO <sub>2</sub> <sup>+</sup> , CO <sub>4</sub> <sup>+</sup> , CO <sup>+</sup> , C <sub>2</sub> O <sub>2</sub> <sup>+</sup> , C <sub>2</sub> O <sub>3</sub> <sup>+</sup> , C <sub>2</sub> O <sub>4</sub> <sup>+</sup> , C <sub>2</sub> <sup>+</sup> , C <sup>+</sup> , CO <sub>3</sub> <sup>-</sup> , CO <sub>4</sub> <sup>-</sup>
O <sub>2</sub> , O, O <sub>3</sub>	$O_2(V_1-V_{15}), O_2(A_1\Delta),$ $O_2(A_3C_3C_1)^{\S}, O_2(B_1\Sigma), O(1D),$ O(1S)	e, O <sup>+</sup> , O <sup>-</sup> , O <sub>2</sub> <sup>-</sup> , O <sub>2</sub> <sup>+</sup> , O <sub>4</sub> <sup>-</sup> , O <sub>4</sub> <sup>+</sup> , O <sub>3</sub> <sup>-</sup>
N <sub>2</sub> , N, NO <sub>2</sub> , NO, N <sub>2</sub> O, N <sub>2</sub> O <sub>3</sub> , N <sub>2</sub> O <sub>4</sub> , N <sub>2</sub> O <sub>5</sub>	$N_2(V_1-V_{21}), N_2(^2D), N_2(^3P),$ $N_2(A_1\Sigma), N_2(A_3\Sigma), N_2(B_3\Pi),$ $N_2(C_3\Pi)$	N <sub>2</sub> <sup>+</sup> , N <sub>3</sub> <sup>+</sup> , N <sub>4</sub> <sup>+</sup> , N <sup>+</sup> , NO <sub>2</sub> <sup>+</sup> , NO <sub>2</sub> <sup>-</sup> , N <sub>2</sub> O <sup>-</sup> , N <sub>2</sub> O <sup>+</sup> , NO <sup>+</sup> , NO <sup>-</sup> , NO <sub>3</sub> <sup>-</sup>
Ar	Ar( ${}^{4}$ S), Ar( ${}^{4}$ P), Ar $_{2}$ (E) ${}^{5}$ , Ar( ${}^{4}$ S $_{1}$ [P $_{0}$ ]), Ar( ${}^{4}$ S $_{1}$ [P $_{1}$ ]), Ar( ${}^{4}$ S $_{1}$ [P $_{2}$ ]), Ar( ${}^{4}$ S $_{1}$ [P $_{1}$ ]),	Ar <sup>+</sup> , Ar <sub>2</sub> <sup>+</sup>

 $<sup>\</sup>S: O_2(A_3C_3C_1)$  is a combination of three electronically excited states  $O_2(A^3\Sigma)$ ,  $O_2(C^3\Delta)$  and  $O_2(c^1\Sigma)$  with a threshold energy of 4.5 eV.

For details about the notations of the other excited levels, see our earlier work  $7^{1,74,103}$ .

<sup>§§:</sup>  $Ar_2(E)$  is a combination of the excited states  $Ar_2(^1\Sigma)$  and  $Ar_2(^3\Sigma)$  of the  $Ar_2$  dimer.