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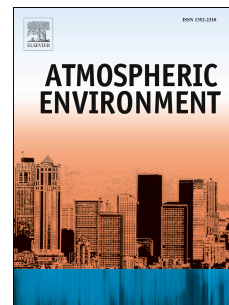
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Quantifying inflow uncertainties in RANS simulations of urban pollutant dispersion

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Abstract

Numerical simulations of flow and pollutant dispersion in urban environments have the potential to support design and policy decisions that could reduce the population's exposure to air pollution. Reynolds-averaged Navier-Stokes simulations are a common modeling technique for urban flow and dispersion, but several sources of uncertainty in the simulations can affect the accuracy of the results. The present study proposes a method to quantify the uncertainty related to variability in the inflow boundary conditions. The method is applied to predict flow and pollutant dispersion in downtown Oklahoma City and the results are compared to field measurements available from the Joint Urban 2003 measurement campaign. Three uncertain parameters that define the inflow profiles for velocity, turbulence kinetic energy and turbulence dissipation are defined: the velocity magnitude and direction, and the terrain roughness length. The uncertain parameter space is defined based on the available measurement data, and a non-intrusive propagation approach that employs 729 simulations is used to quantify the uncertainty in the simulation output. A variance based sensitivity analysis is performed to identify the most influential uncertain parameters, and it is shown that the predicted tracer concentrations are influenced by all three uncertain variables. Subsequently, we specify different probability distributions for the uncertain inflow

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variables based on the available measurement data and calculate the corresponding means and 95% confidence intervals for comparison with the field measurements at 35 locations in downtown Oklahoma City.

Keywords: urban flow, dispersion, Reynolds-Averaged Navier-Stokes, uncertainty quantification, Joint Urban 2003

1. Introduction

2 The latest World Health Organization (WHO) factsheet dedicated to ambient
3 pollution shows that 92% of the world population lives in areas where
4 the WHO air quality guidelines are not met [1]. The capability to provide
5 accurate predictions of the dispersion of harmful pollutants in urban areas
6 would enable supporting design and policy decisions to alleviate this problem.
7 This motivates research into the use of computational fluid dynamics (CFD)
8 to predict flow and dispersion phenomena in complex urban areas. A vast
9 review of CFD studies of atmospheric boundary layer (ABL) flows in urban
10 environments was recently published by Blocken [2]. The paper considered
11 both Reynolds-averaged Navier-Stokes (RANS) [3, 4, 5, 6] and large-eddy
12 simulations (LES) [7, 8, 9, 10, 11]. LES has the advantage of providing a more
13 accurate representation of the turbulent flow features, since the larger, energy
14 containing scales are resolved and only the small scales are modeled. Its
15 main drawback is the significant amount of computational resources required,
16 which is still limiting its use in complex engineering applications. As a result,
17 more computationally affordable RANS simulations, which model the entire
18 spectrum of turbulent scales, are commonly used in the wind engineering
19 practice [12]. Several previous studies have performed steady and unsteady
20 RANS simulations of dispersion of full-scale experimental campaigns [13, 14].
21 These studies focused on evaluating model performance by running a set of
22 simulations with deterministic inflow conditions, and comparing their results
23 with experimental data. Traditional performance metrics, such as fractional
24 bias and normalized mean square error [15], were calculated, and varying
25 degrees of success were reported, depending on the experimental period of
26 time, the atmospheric stratification, and the intensity and variability of the
27 inflow winds. These results indicate that the accuracy of RANS simulation
28 results for urban flow and dispersion can be compromised by several sources
29 of uncertainty.

30

31 Traditionally, a distinction is made between two types of uncertainty:
32 the epistemic uncertainty, which is related to a lack of knowledge, and the
33 aleatory uncertainty, which is intrinsic to the nature of the system and can
34 not be reduced. In RANS simulations, the epistemic uncertainty largely
35 originates from the turbulence modeling, and the development of methods to
36 quantify this uncertainty is a relatively recent area of research [16, 17, 18, 19].
37 An important aleatory uncertainty is the definition of the inflow boundary
38 conditions. These are influenced by the larger-scale atmospheric conditions
39 and can therefore be highly variable [20, 21].

40 Previous research focused on quantifying the latter uncertainty in sim-
41 ulations of the flow in downtown Oklahoma City [22, 23]. In these studies
42 we defined three uncertain parameters that define the inflow boundary con-
43 dition: the wind speed and direction, and the ABL roughness length. To
44 fully evaluate the capabilities of the approach we considered a test case for
45 which field experiments are available for comparison to the UQ results. We
46 selected the Joint Urban 2003 (JU2003) measurement campaign, which pro-
47 vides a full month of measurements in downtown Oklahoma City. We focused
48 specifically on 30-minute time series of wind flow measurements at 13 loca-
49 tions in the downtown area using portable wind detectors (PWIDS). During
50 the measurement campaign, the predominant wind direction was from the
51 south, thus wind measurements from the south-most station were used to
52 define probability density functions (pdfs) for the uncertain inflow param-
53 eters. A non-intrusive uncertainty quantification (UQ) approach based on
54 729 deterministic RANS simulations was used to propagate the inflow un-
55 certainties to the simulation output. The resulting wind predictions for the
56 mean and 95% confidence intervals were compared to the field measurements,
57 showing that the 95% UQ confidence intervals encompassed all stations for
58 the experimental mean velocity magnitude and 11 out of 13 stations for the
59 experimental mean velocity direction.

60 The objective of the present research is to extend this study to quan-
61 tify the inflow uncertainty in simulations of urban dispersion, and to include
62 a more detailed analysis of the methodology and results. We focus on a
63 different measurement period from the JU2003 campaign, considering one
64 of the intensive observation periods that included a continuous tracer re-
65 lease to analyze dispersion. Wind and concentration measurements are com-
66 pared to the UQ results using 35 portable and super portable wind detectors
67 (PWIDS/SPWIDS), and 8 real time tracer analyzers located in the down-
68 town area. As in [22, 23] we define three uncertain parameters to define the

69 inflow boundary conditions and we use a non-intrusive propagation method.

70 As a first step we perform a variance-based sensitivity analysis using the
71 Sobol indices, which quantify the direct and combined contribution of the
72 different uncertain inflow parameters to the variance of the quantities of in-
73 terest. This type of analysis enables identifying the influential uncertain
74 variables, and the conclusions could potentially reduce the number of evalu-
75 ations needed for future UQ studies in urban environments. Subsequently we
76 investigate the influence of the choice for the input pdfs by considering two
77 different beta distributions fitted to the experimental data. The first repre-
78 sents the full variance of the experimental signal, and therefore represents
79 all fluctuations in the measured signal as an inflow uncertainty. The second
80 beta distribution represents a reduced variance, where the variance corre-
81 sponding to the turbulence kinetic energy in the RANS model is removed
82 from the measured signal, and only the remaining portion of the measured
83 variance is considered an inflow uncertainty. Mean and 95% confidence inter-
84 vals for the corresponding flow and tracer concentration fields are presented
85 and compared to the available experimental data.

86 The remainder of this paper is organized in six sections. Section 2 summa-
87 rizes the Joint Urban 2003 (JU2003) field measurements. Section 3 provides
88 the details of the computational model, including the governing equations
89 and their discretization, the boundary conditions applied, and the computa-
90 tional domain and mesh. Section 4 presents the uncertainty quantification
91 (UQ) methodology, discussing the propagation algorithm, the calculation of
92 the Sobol indices and the definition of the uncertain inflow parameters based
93 on the experimental data. The sensitivity analysis and the comparison be-
94 tween field measurements and statistics for the flow and tracer concentration
95 are presented in section 5. Conclusions and future research directions are
96 formulated in section 6.

97 **2. Experimental Campaign**

98 The JU2003 experiments were performed over one full month during the
99 summer of 2003 and provide day- and nighttime data for different boundary
100 layer stratifications [24]. Previous modeling studies have used the database
101 for evaluating models with differing degrees of complexity, and considering
102 different scales. Chan and Leach [25] used it to evaluate and further develop
103 the unsteady 3D finite element CFD code FEM3MP, while Neophytoua et
104 al. [26] compared the measurements to results of different numerical ap-

105 proaches, ranging from rather simple to complex methods. In both studies
 106 the authors highlight the importance of the inflow boundary conditions to
 107 produce realistic urban flow predictions.

108 The specific data used for this study was obtained during one of the inten-
 109 sive observation periods (IOP9) with near-neutral flow conditions when sulfur
 110 hexafluoride (SF_6) was released in a continuous manner [27]. Specifically, the
 111 NOAA Air Resources Laboratory Field Research Division operated 8 mobile
 112 real-time analysers (ARLFRD) that provide concentration measurements.
 113 Wind measurements were performed by Dugway Proving Ground systems
 114 using 20 super portable wind detectors (SuperPWIDS) and 15 portable wind
 115 detectors (PWIDS), which are sonic and 3D sonic anemometers with sam-
 116 pling frequencies of 10Hz and 1Hz respectively. Figure 1 displays the location
 117 of the different measurement stations in downtown Oklahoma city, as well as
 118 the SF_6 source location. PWIDS 15, the south-most sensor, is used to define
 119 the uncertainty in the inflow variables as described in Section 4. All other
 120 sensor measurements will be used for validation of the model results.

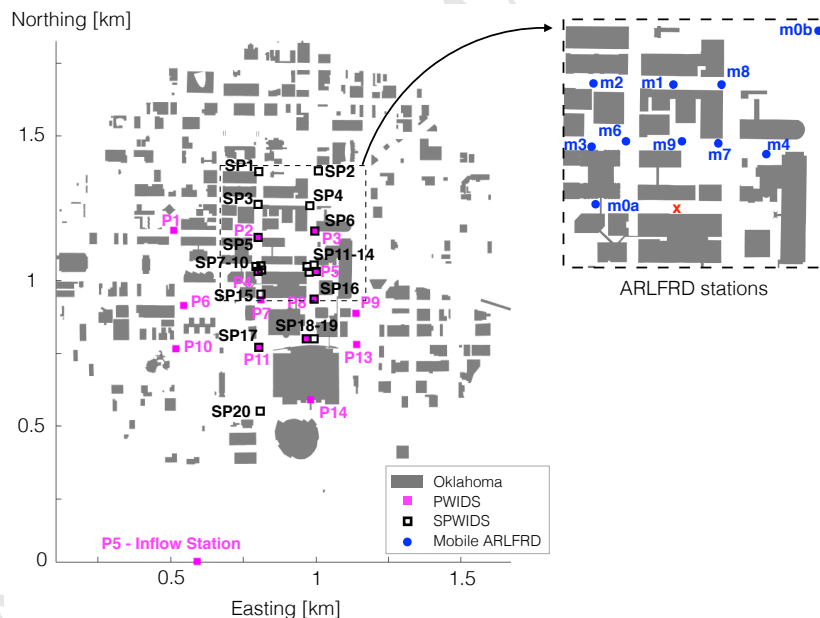


Figure 1: JU2003 measurement stations: portable wind detector (P, magenta squares), super portable wind detectors (SP, black squares) and tracer measurement mobile real time analysers (m, blue circles). SF_6 release location (red cross).

121 3. Computational model

122 3.1. Governing equations, discretization and solution method

123 The governing equations for the flow field are identical to those used in
 124 [22, 23]. The steady, incompressible Reynolds-averaged Navier-Stokes equa-
 125 tions are used to solve the neutrally stratified ABL flow. Turbulence closure
 126 is obtained using the standard two equation k- ε turbulence model, which
 127 computes the Reynolds stresses using the linear eddy viscosity hypothesis
 128 [28].

129 To predict the concentration field, an additional transport equation for a
 130 passive scalar is solved [29]:

$$\frac{\overline{u}_j}{\partial x_j} \frac{\partial \overline{C}}{\partial x_j} - \frac{\nu_t}{Sc_t} \frac{\partial^2 \overline{C}}{\partial x_j \partial x_j} = C_{source} \quad (1)$$

131 where \overline{u}_j are the time-averaged velocity components, and \overline{C} represents the
 132 time-averaged tracer concentration. Turbulent dispersion is represented us-
 133 ing a standard gradient diffusion model, where the turbulent diffusion coef-
 134 ficient is defined as the ratio of the turbulent viscosity, ν_t , and the turbulent
 135 Schmidt number, Sc_t , which is fixed to a constant value of 0.7. The tracer
 136 is introduced continuously in the simulation in the cell corresponding to the
 137 location of the release of interest during IOP9 [24], as shown in Figure 1.

138 The equations were solved numerically using the open source set of li-
 139 braries, OpenFOAM [30]. The flow equations are discretized using the finite-
 140 volume method, with second order upwind schemes for velocity, turbulence
 141 kinetic energy and turbulence dissipation. The simulations were run for 3000
 142 iterations, which resulted in a converged solution with residuals below 1e-06
 143 for all variables.

144 3.2. Computational domain and boundary conditions

145 The domain used to run the CFD simulations is presented in figure 2.
 146 The size of the domain was determined following the best practice guidelines
 147 from COST action 732 [31], which lead to a domain size of 4km by 3.9km
 148 in horizontal direction. The vertical domain height is 1.2 km, which equals
 149 $8H_{max}$, where $H_{max} = 150\text{m}$ is the height of the tallest building in downtown
 150 Oklahoma City. All lateral boundaries are also located at a distance larger
 151 than $5H_{max}$ from the nearest building in order to avoid any unwanted effects
 152 from the boundary conditions on the flow solution. This design allows using

153 all lateral boundaries as either inlets or outlets, thereby efficiently enabling
 154 simulations for the different wind directions (70° to 120°) that occurred dur-
 155 ing IOP9. More specifically, a 90° wind direction, which corresponds to the
 156 wind coming from the south, would use the purple boundary from Figure
 157 2 as an inflow, and the white boundary as an outflow. For inflow direc-
 158 tions smaller than 90° , the blue boundary is an additional inflow, while the
 159 pink boundary is an additional outlet. The opposite would hold for inflow
 160 directions larger than 90° .

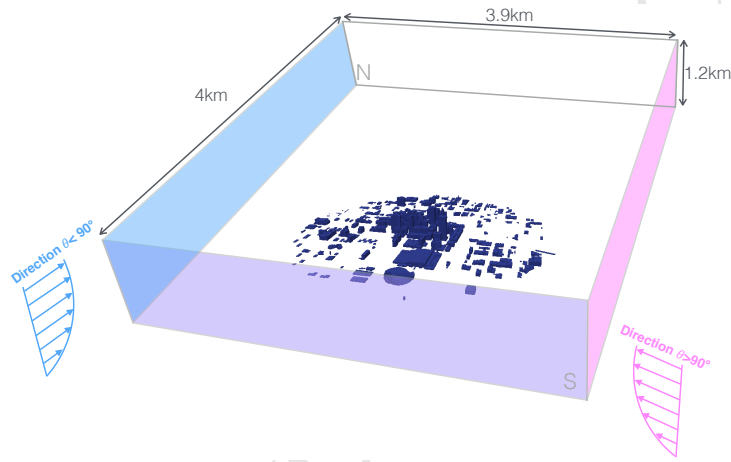


Figure 2: Computational Domain (N: north, S: south, South-North direction represents 90 degrees orientation).

161 At the outlets zero-gradient boundary conditions are imposed. On the
 162 inlets, the inflow boundary conditions represent a fully developed neutral
 163 surface layer [32], since the time-period selected from the experiments has
 164 near-neutral conditions. The surface layer profiles were used over the entire
 165 height of the domain. The influence of the profiles specified above 100m is
 166 expected to be negligible when considering quantities of interest at pedestrian
 167 level. The boundary conditions thus specify a logarithmic profile for velocity,
 168 a constant profile for turbulence kinetic energy, and a profile for epsilon
 169 that assumes equilibrium between turbulence kinetic energy production and
 170 dissipation:

$$U = \frac{u_*}{\kappa} \ln \left(\frac{z + z_0}{z_0} \right) \quad (2)$$

171

$$k = \frac{u_*^2}{\sqrt{C_\mu}} \quad (3)$$

172

$$\varepsilon = \frac{u_*^3}{\kappa(z + z_0)} \quad (4)$$

173 U represents the mean velocity, κ the von Karman constant (0.41), k and ε
 174 the turbulence kinetic energy and dissipation rate, and C_μ is a turbulence
 175 model constant (0.09). The friction velocity is determined from the specified
 176 aerodynamic roughness, and the reference velocity, U_{inflow} , measured at 30m
 177 height by PWIDS 15 (figure 1).

178 At the top boundary a constant shear stress boundary condition is mim-
 179 icked by specifying the velocity, turbulence kinetic energy and turbulence
 180 dissipation rate as the values given by the inlet profiles at the height of the
 181 top boundary. At wall boundaries a no-slip condition is imposed and wall
 182 functions are used to calculate the flow variables in the wall-adjacent cells. At
 183 the ground boundary an ABL wall function based on a constant ABL rough-
 184 ness length, z_0 , is used. This wall function ensures that the neutral ABL
 185 inflow profiles for velocity, turbulence kinetic energy and turbulence dissipa-
 186 tion rate are maintained downstream of the inlet [33, 34]. At the building
 187 walls a standard log law wall function for smooth surfaces is applied.

188 3.3. Computational mesh and grid dependency

189 The parallel mesh generator, SnappyHexMesh [30], is used to design the
 190 computational mesh. It consists of ~ 7.4 million cells, the majority of which
 191 are hexahedral. The mesh is progressively refined to achieve the highest res-
 192 olution in the downtown area where the tracer concentration measurements
 193 were performed. In this region, the vertical resolution ranges from 0.69 to 2m
 194 near the ground, while the near wall resolution on the building walls varies
 195 from 0.3 to 1.32m. Before performing the UQ study, a grid dependency study
 196 was performed by running simulations with a finer mesh consisting of 16.8
 197 million cells.

198 The inflow parameters for the grid dependency study are specified to
 199 match the time-average of the measurement at PWIDS 15 during IOP9, be-
 200 ing $U = 6.13m s^{-1}$, $z_0 = 0.29m$, and the wind direction $\theta = 97^\circ$. Figure 3
 201 presents a comparison of the velocity at the PWIDS stations between the
 202 time-averaged measurements and the RANS results from the two different
 203 meshes. For both cases we computed hit rates with the experimental data

204 [31]. The thresholds to produce a 'hit' for the velocity magnitude were 0.25
 205 for the relative discrepancy, and 0.008 for the absolute values. For the ve-
 206 locity direction only an absolute threshold of 20° was used. We found no
 207 improvement when using the finer mesh compared to the nominal one. In
 208 addition, we calculated hit rates for the nominal mesh using the fine mesh
 209 results as the observation values, which produced a hit in more than 75 % of
 210 the measurement stations for both wind magnitude and direction. A similar
 211 analysis was performed for the predicted pollutant concentrations, which also
 212 showed no improvements when using the fine mesh compared to the nominal
 213 mesh.

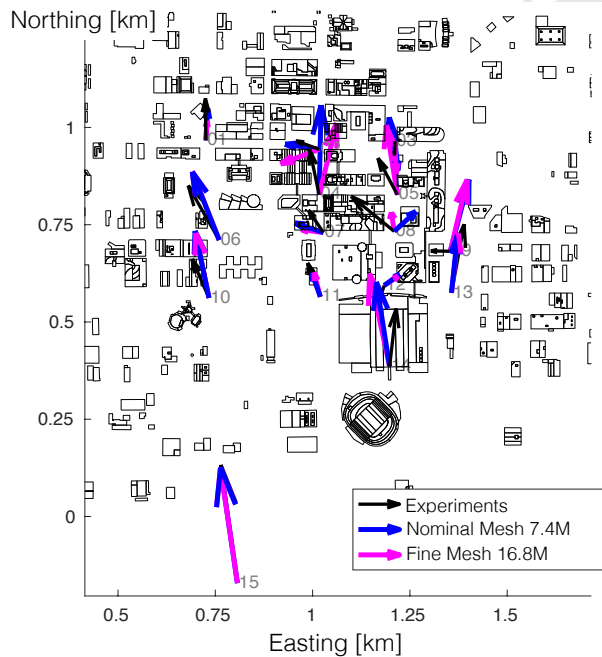


Figure 3: Comparison of the velocity vectors at the PWIDS stations: time-averaged measurements, RANS nominal mesh, and RANS fine mesh results.

214 Considering the relatively small changes in the predictions obtained with
 215 the nominal and fine mesh, and given the large number of simulations re-
 216 quired for the UQ analysis, the nominal mesh has been employed for the
 217 remainder of the study.

218 4. Inflow Uncertainty Quantification Method

219 We characterize the uncertainty in the inflow boundary conditions spec-
220 ified by Eqs. 2-4 using three uncertain parameters: the wind direction, the
221 wind magnitude at a reference height, and the upstream terrain roughness
222 height z_0 . The wind magnitude at a reference height is used to calculate
223 the friction velocity u_* , thereby influencing the profiles for U , k and ϵ . In
224 the following sections we first address the definition of probability distribu-
225 tions for the three uncertain parameters, and subsequently we summarize
226 the UQ method, the calculation of the Sobol indices and the mean and 95%
227 confidence intervals for the quantities of interest (QoI).

228 4.1. Characterization of uncertain inflow parameters

229 Characterizing the uncertain input parameters is an essential component
230 of a UQ framework. For the initial sensitivity analysis we assumed uniform
231 distributions for all three variables. The minima and maxima for the veloc-
232 ity magnitude and direction were obtained from the field measurements at
233 PWIDS 15. This station is selected because it is the south-most measurement
234 location and the wind direction during the measurements was predominantly
235 from the south. The range for the aerodynamic roughness height was based
236 on a roughness map of Oklahoma City available from [35].

237 For calculating the means and 95% confidence intervals, we propose a
238 method to define more realistic, non-uniform, probability distributions. For
239 the ABL roughness height, the available roughness map was digitally pro-
240 cessed considering the area south of the city to build an approximate dis-
241 tribution for z_0 , as presented in figure 4. Probability distributions for the
242 velocity magnitude and direction are obtained from the experimental time
243 series measured at PWIDS 15. Figure 4 presents the distributions from the
244 30 minute period of interest corresponding to the first continuous release
245 during IOP9.

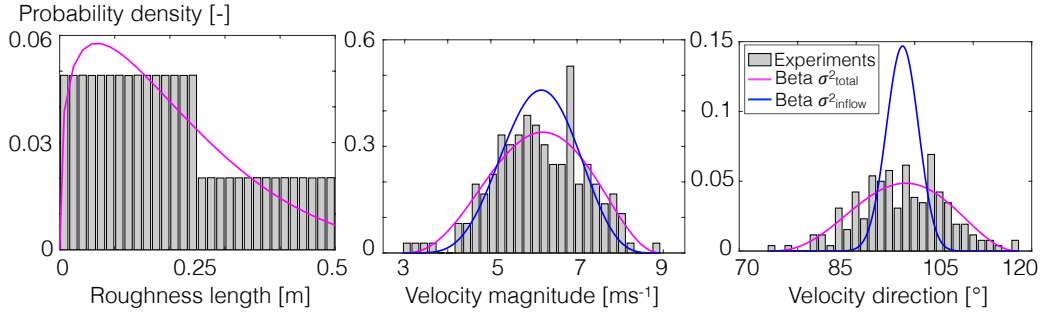


Figure 4: Probability density functions for the inflow uncertain parameters: roughness length, velocity magnitude and direction.

246 Figure 4 also depicts the beta probability density functions that were
 247 fitted to these distributions. These will be used to generate realizations from
 248 the spectral expansions for the quantities of interest, such that the mean
 249 and variance can be computed. For the velocity magnitude and direction,
 250 two different beta distributions will be considered. The coefficients for the
 251 first one, beta σ_{total}^2 , were determined to match the mean and the variance
 252 of the experimentally measured time-series. The second one, beta σ_{inflow}^2 ,
 253 also matches the mean of the experiment but has a reduced variance. The
 254 rationale behind this second approach is that the turbulence model employed
 255 in the RANS simulations already accounts for part of the fluctuations in the
 256 measurements, hence this portion of the variance should not be represented
 257 by variability in the inflow conditions.

258 This 'variance reduction' methodology is based on the approach proposed
 259 by Vervecken et al. [36]. The conceptual idea is that the wind direction vari-
 260 ance can be decoupled into two different contributions: a 'model' component
 261 that is represented by the turbulence model in the RANS simulation, and
 262 an 'external' or 'inflow' component related to the larger-scale variations in
 263 the atmospheric boundary layer, which are not represented by the RANS
 264 turbulence model. The fluctuating velocity is written as the sum of two
 265 different contributions: (1) the fluctuations represented by the RANS tur-
 266 bulence model, u'_m and v'_m ; and (2) the fluctuations induced by the 'exter-
 267 nal' or 'inflow' variability, u'_e and v'_e . It is assumed that the instantaneous
 268 spanwise velocity is much smaller than the instantaneous streamwise veloci-
 269 ty ($v \ll u$), and that the fluctuations are small ($u' \ll u$). In addition
 270 the covariance of the modeled and external fluctuations is neglected. Using

271 these assumptions, the variance in the wind direction can be decomposed as
 272 follows:

$$\sigma_{\theta}^2 \approx \sigma_{\theta, inflow}^2 + \sigma_{\theta, m}^2, \quad (5)$$

273 where the modeled component can be approximated using

$$\sigma_{\theta, m}^2 \approx \frac{2/3k}{U^2}. \quad (6)$$

274 In previous works [23], we applied this relationship to define a beta pdf
 275 with reduced variance $\sigma_{\theta, inflow}^2$ for the wind direction. In the present work
 276 we extend this approach to also define a beta distribution with a reduced
 277 variance for the wind magnitude. Using the same assumptions as above, the
 278 total variance for the velocity magnitude can be written as:

$$\sigma_U^2 \approx \sigma_{U, inflow}^2 + \sigma_{U, m}^2, \quad (7)$$

279 where the modeled contribution can be approximated as:

$$\sigma_{U, m}^2 \approx \frac{2}{3}k \quad (8)$$

280 The coefficients for the reduced variance beta distributions are therefore ob-
 281 tained by conserving the original mean from the PWIDS 15 measurement,
 282 and using the variances $\sigma_{\theta, inflow}^2$ and $\sigma_{U, inflow}^2$.

283 This variance reduction can be interpreted as a filtering operation on the
 284 original measurement data. If we assume the smaller scale fluctuations are
 285 represented by the turbulence model in the RANS simulations, we can find
 286 the cut-off frequency that accounts for the portion of the variance repre-
 287 sented by $\sigma_{\theta, m}^2$ and $\sigma_{U, m}^2$. Figure 5 presents the normalized spectral energy
 288 density, indicating the cut-off frequency that separates the larger scale inflow
 289 variations from the smaller scale fluctuations represented by the turbulence
 290 model. Figure 6 presents the corresponding original and filtered signals for
 291 the velocity magnitude and direction at PWIDS 15. In section 5.2 we com-
 292 pare results for the means and 95% confidence intervals of the velocity and
 293 concentration field using both distributions to demonstrate the importance
 294 of the choice of the uncertain parameter space.

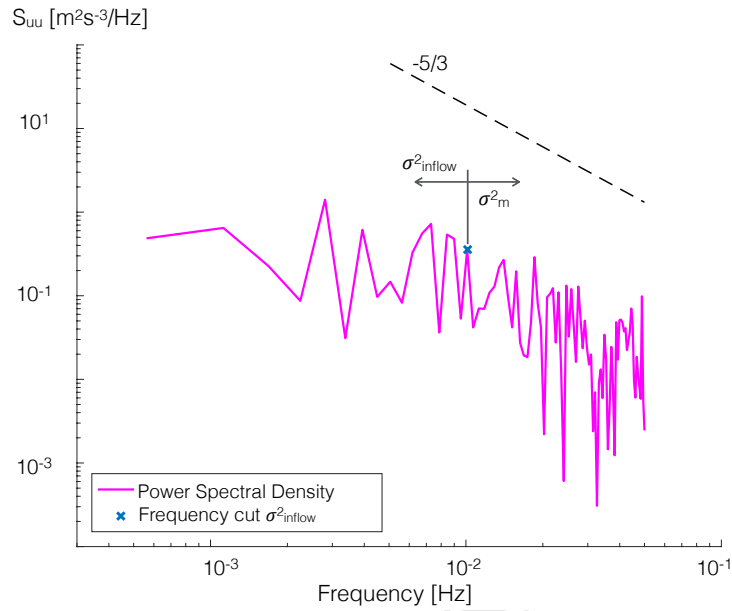


Figure 5: Normalized spectral energy density in function of frequency at PWIDS 15 station.

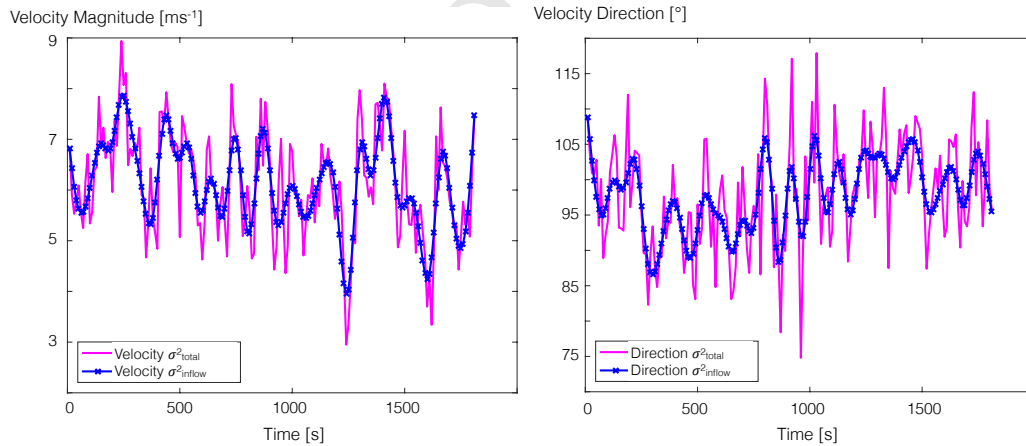


Figure 6: Original velocity magnitude and direction signals compared to the filtered signals considering uniquely inflow variance.

295 4.2. Propagation algorithm

296 The non-intrusive UQ code RAY developed at the Stanford UQ lab [37]
 297 was employed to propagate the inflow uncertainties to the output quanti-

ties of interest. The code requests user input regarding (1) the uncertain parameter space, i.e. minimum and maximum values for each uncertain parameter, (2) the quadrature method to be used, and (3) the maximum allowed number of simulations to perform the analysis. Using this information RAY determines the collocation points, i.e. the points in the 3D uncertain input parameter space for which simulations will be performed. Using the output of these simulations RAY constructs a spectral expansion and calculates statistics and probability distributions of the quantities of interest. The following sections subsequently describe the construction of the spectral expansion, the calculation of the Sobol indices, and the calculation of the statistics of the QoI in more detail [38].

4.2.1. Constructing the spectral expansion

The three uncertain inflow parameters, ξ_j are assumed to be mutually independent and uniform random variables. The n quantities of interest u^k are a function of ξ_j :

$$u^K = u^k(\xi_1, \xi_2, \xi_3) \quad \text{with} \quad k = 1, \dots, n \quad (9)$$

Each variable ξ_j is defined on an interval Ξ_j , and the uncertain parameter space is characterized by the corresponding hypercube $\Xi = \Xi_1 \times \Xi_2 \times \Xi_3$. A probabilistic characterization of u^K can be obtained by formulating a spectral expansion:

$$u^K = \sum_{k=0}^{\infty} u_i^k \psi_i(\xi_1, \xi_2, \xi_3) \approx \sum_{i=0}^P u_i^k \psi_i(\xi_1, \xi_2, \xi_3), \quad (10)$$

where P is the order of truncation of the series and ψ_i are multivariate orthogonal polynomials. In the present study we employ Legendre polynomials and the maximum order of the polynomial terms is 4. The polynomial coefficients u_i^k are evaluated non-intrusively using Galerkin projection, with the integrals computed using Clenshaw-Curtis quadrature on an isotropic tensor grid:

$$u_i^k = \frac{\int_{\Xi} u^k \psi_i d\xi_1 d\xi_2 d\xi_3}{\int_{\Xi} \psi_i \psi_i d\xi_1 d\xi_2 d\xi_3} = \frac{\sum_{l=0}^N u^k(\xi_1^l, \xi_2^l, \xi_3^l) \psi_i(\xi_1^l, \xi_2^l, \xi_3^l) w^l}{\sum_{l=0}^N \psi_i(\xi_1^l, \xi_2^l, \xi_3^l) \psi_i(\xi_1^l, \xi_2^l, \xi_3^l) w^l}, \quad (11)$$

where $(\xi_1^l, \xi_2^l, \xi_3^l)$ is point l of the isotropic tensor grid (abscissa) and w^l are the integration weights.

325 *4.2.2. Calculating the Sobol Indices*

326 Direct and combined Sobol indices are commonly used to analyse the
 327 importance of the different uncertain parameters in a UQ analysis [39, 40].
 328 The indices quantify the relative contribution of each uncertain variable to
 329 the variance of a QoI. The calculation of the Sobol indices is straightfor-
 330 ward when the uncertain variables have probability density functions that
 331 are orthogonal to the polynomial basis. Hence, for this sensitivity analysis,
 332 we assume uniform distributions for all uncertain variables. The direct and
 333 combined Sobol indices can then be calculated as follows:

$$S_i = \frac{\sigma^2 [\mathbb{E} (u|\xi_i)]}{\sigma^2(u)} \quad (12)$$

334

$$S_{ij} = \frac{\sigma^2 [\mathbb{E} (u|\xi_i, \xi_j)]}{\sigma^2(u)} \quad (13)$$

335 The direct indices, S_i , provide a measure of the contribution of an isolated
 336 uncertain parameter, ξ_i , to the variance of the QoI. The combined indices,
 337 S_{ij} , quantify the contribution of a combination between different uncertain
 338 parameters ξ_i and ξ_j . The sum of all direct and combined Sobol indices is
 339 equal to 1. The expectations and variances required for this analysis are
 340 also easily computed by exploiting the orthogonality of the densities and the
 341 polynomial basis:

$$\mathbb{E} [u^K] = \int_{\Xi} u^k d\xi_1 d\xi_2 d\xi_3 \approx \sum_{i=0}^P u_i \Psi_i(\xi) \approx u_0 \quad (14)$$

342 while the variance would be calculated following:

$$\sigma^2 [u^K] = \mathbb{E} [(u^K - \mathbb{E} [u^K])^2] \approx \sum_{i=1}^P u_i^2 \langle \Psi_i(\xi) \rangle^2 \quad (15)$$

343 *4.2.3. Calculating statistics for the output quantities of interest*

344 For uncertain input variables with beta distributions as presented in fig-
 345 ure 4 a different approach is required to compute statistics for the flow and
 346 concentration field. Once the spectral response surfaces for the quantities of
 347 interest are determined, probability density functions for u^K can be obtained
 348 by drawing realizations from the spectral representations according to the

349 beta distributions for the uncertain input variables. The mean and 95% con-
350 fidence intervals are subsequently determined from these probability density
351 functions. In the present analysis 10.000 realizations were used to calculate
352 the statistics.

353 5. Results

354 In this section, we first present the sensitivity analysis. Subsequently, the
355 analysis of the results focuses on the predicted mean and 95% confidence
356 intervals (CI) for wind magnitude, direction and non-dimensional concentra-
357 tion.

358 5.1. Variance based sensitivity analysis

359 The analysis of the Sobol indices, which are calculated as explained in
360 section 4, aims to investigate which of the uncertain input parameters con-
361 tributes more to the variance of the quantities of interest. In the following
362 the Sobol indices are presented together with the corresponding variance,
363 since the analysis is particularly relevant in regions of the flow where the
364 variance is larger. Figures 7, 8 and 9 therefore first present a contour plot for
365 the variance of the quantity of interest under consideration in the upper left.
366 Subsequently, from left to right and top to bottom, the other contour plots
367 show the corresponding direct Sobol indices for the three uncertain variables
368 (S_U , S_θ , and S_{z_0}), and two combined Sobol indices ($S_{U\theta}$ and $S_{\theta z_0}$).

369 5.1.1. Wind flows

370 Figure 7 presents the variance of the non dimensional velocity magni-
371 tude in downtown Oklahoma City and the corresponding direct and two of
372 the combined Sobol indices. The variance contour plot for the predicted
373 non-dimensional velocity magnitude shows that the variance becomes non
374 negligible after the first row of buildings, and therefore the investigation of
375 the Sobol indices is most relevant in those locations.

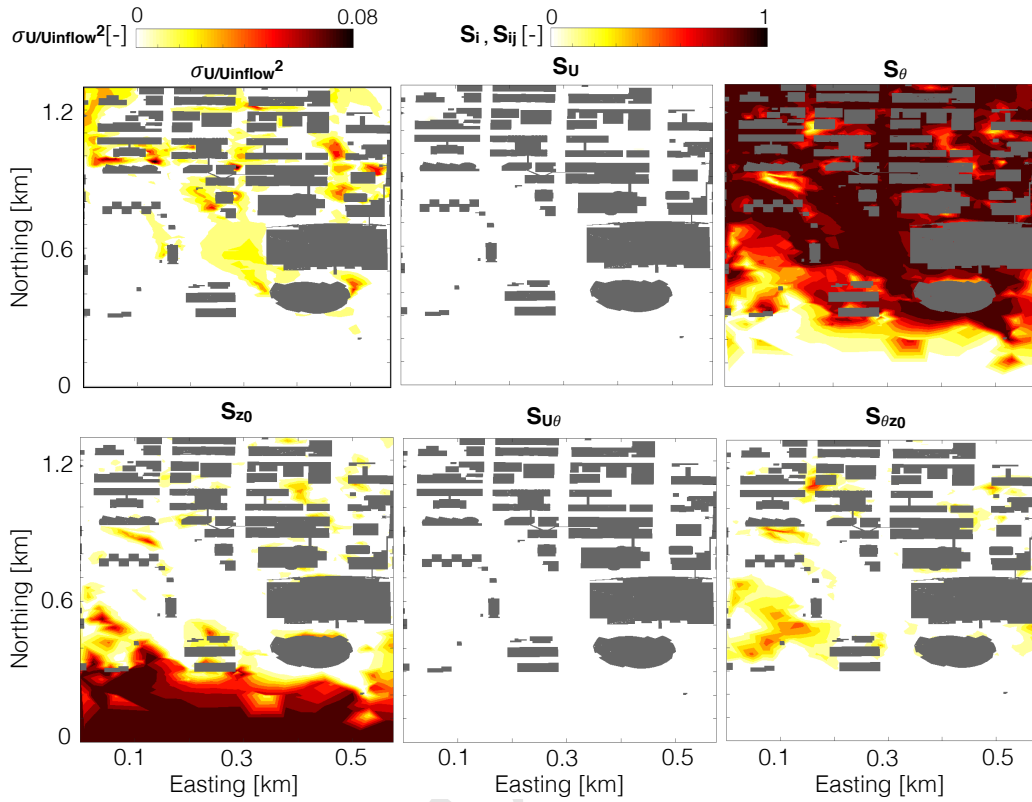


Figure 7: Sensitivity analysis for the prediction of the non-dimensional velocity magnitude U/U_{inflow} in downtown Oklahoma City: contour plots for the variance and the direct and combined Sobol indices.

376 The top center plot shows how the contribution of the uncertainty in the
 377 inflow velocity magnitude, quantified by S_U , is negligible. Downstream of
 378 the first row of buildings the most important contribution to the variance
 379 originates from uncertainty in the inflow wind direction as shown in the plot
 380 for S_θ . Upstream of the city the roughness length has the largest contribution
 381 to the variance, but the magnitude of the variance is small in this region.
 382 Some isolated locations in the flow field show additional sensitivity to the
 383 roughness length or to the combined effect between the inflow wind direction
 384 and the roughness length. These areas seem to correspond to larger recircu-
 385 lation zones. Combined sobol indices from the velocity magnitude with the
 386 other parameters shown negligible effects on the quantity of interest.

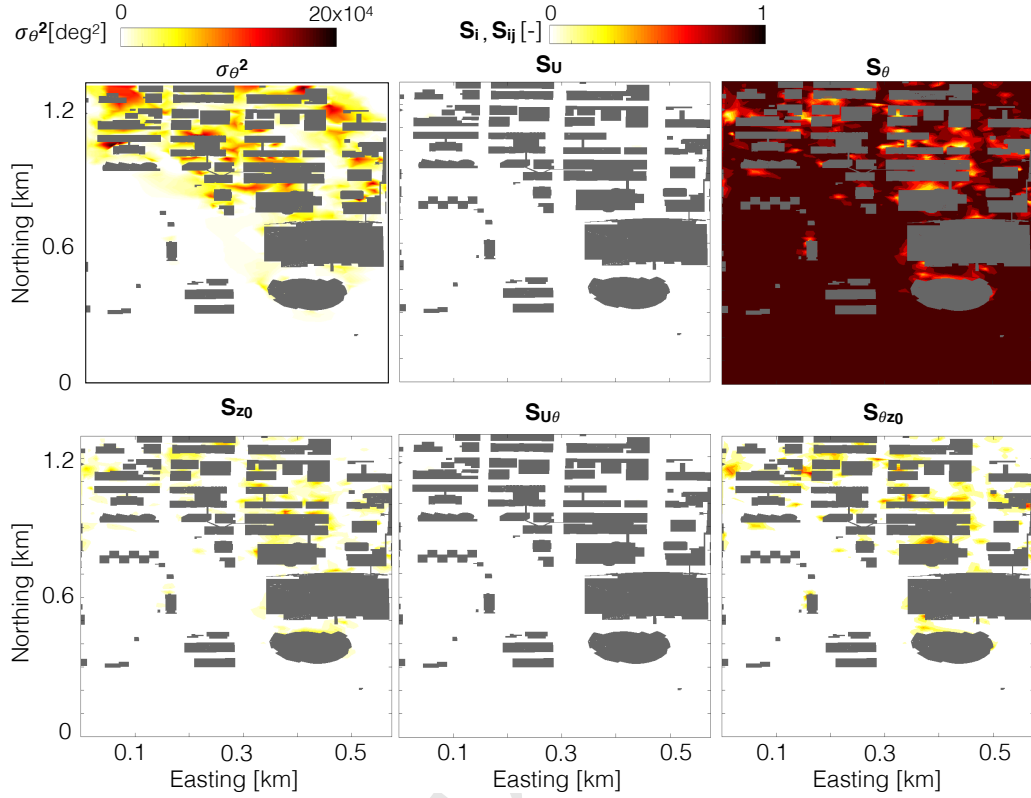


Figure 8: Sensitivity analysis for the prediction of the velocity direction θ in downtown Oklahoma City: contour plots for the variance and the direct and combined Sobol indices.

387 Figure 8 presents the variance of the flow direction in downtown Okla-
 388 homa City and the corresponding direct and combined Sobol indices. As
 389 for the velocity magnitude, the contour plot of the variance shows that the
 390 variance becomes non negligible after the first row of buildings. The con-
 391 tour plots of the Sobol indices clearly identify the uncertainty in the inflow
 392 velocity direction as the dominant factor.

393 5.1.2. Dispersion

394 Figure 9 depicts the variance of the non-dimensional tracer concentration
 395 and the corresponding Sobol indices. The non-dimensional concentration is
 396 defined as:

$$K = \frac{C^* H^2 U_{ref}}{Q_S} \quad (16)$$

397 where C^* represents the ratio between the local and the source concentra-
 398 tions, U_{ref} is the inflow velocity at 2m height, H is the average building
 399 height, set to 30m, and Q_S is the source flow rate.

400 The plots demonstrate that all three uncertain inflow parameters have a
 401 non-negligible effect. Near the source the variance is highest and the wind di-
 402 rection is shown to be the most influential parameter. As one would expect, it
 403 strongly determines the initial plume trajectory. Further downstream of the
 404 source, the sobol index for the wind direction decreases, and the inflow ve-
 405 locity magnitude and roughness length have a significant contribution to the
 406 variance. The combined Sobol indices for the wind magnitude and direction,
 407 and for the wind direction and roughness length also show non negligible con-
 408 tributions, indicating a complex interaction between the uncertain variables
 409 and the quantity of interest.

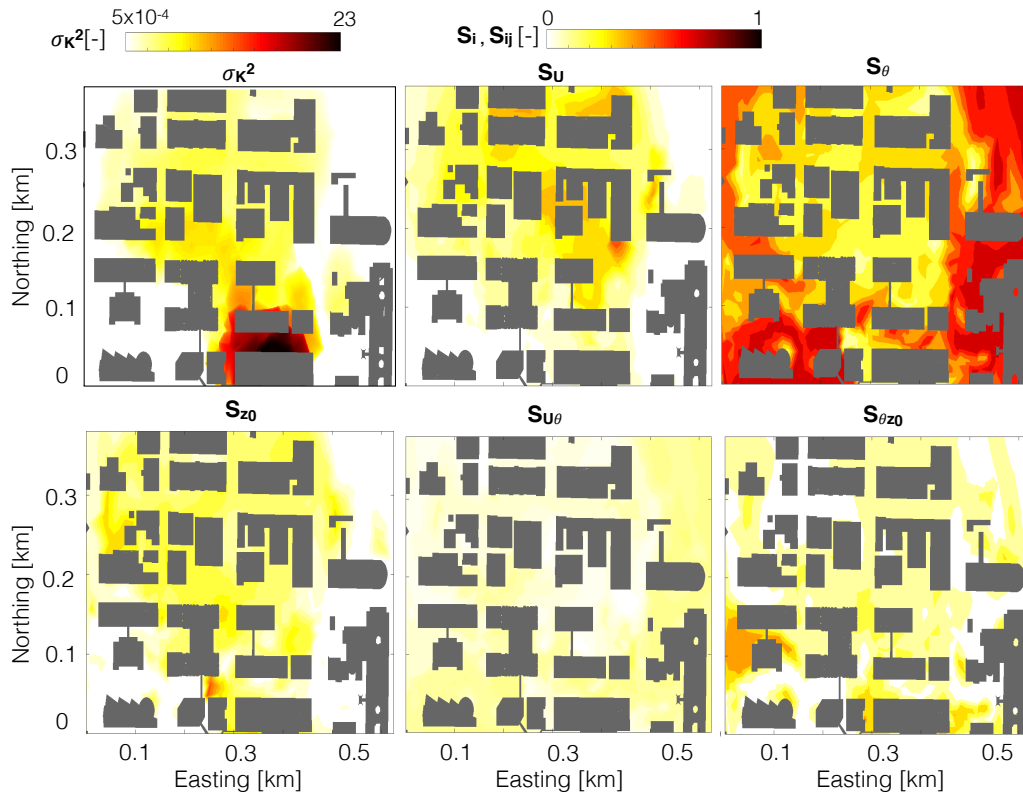


Figure 9: Sensitivity analysis for the prediction of the non-dimensional concentration K in downtown Oklahoma City: contour plots for the variance and the direct and combined Sobol indices.

410 *5.2. Comparison of the predicted mean and 95% Confidence Intervals*

411 The main objective of using the proposed UQ framework is to add con-
412 fidence intervals to the simulation results that reflect the uncertainty in the
413 solution related to variability in the inflow conditions. The confidence in-
414 tervals should encompass the time-averaged field data measured during the
415 period of interest, provided that other uncertainties are sufficiently small.
416 The results for the wind field are presented by plotting the time-averaged
417 experimental velocity vector at the measurement stations and compare that
418 to the mean and 95% CI for the velocity predicted by the UQ study. The
419 results for the concentration field are presented by plotting the mean with
420 95% CI bars for each measurement station.

421 *5.2.1. Wind flow results*

422 Figure 10 presents the results for the velocity field at the 13 PWIDS sta-
423 tions in the downtown area. The black arrows represent the velocity vectors
424 from the field measurements, time-averaged over the experimental period of
425 interest. The magenta and blue arrows show the mean from the UQ study,
426 respectively obtained by assuming the previously presented beta distribu-
427 tions with the total and reduced variance for the inflow wind magnitude and
428 direction (see Figure 4). The corresponding 95% confidence intervals for the
429 predicted velocities are defined as the values where the cumulative density
430 function is 0.025 and 0.975. They are represented in the plot by sectors of an
431 annulus, where the 95% CI of the velocity magnitude is represented by the
432 radii of the annulus, while the 95% CI for the wind direction is represented
433 by the sector angle.

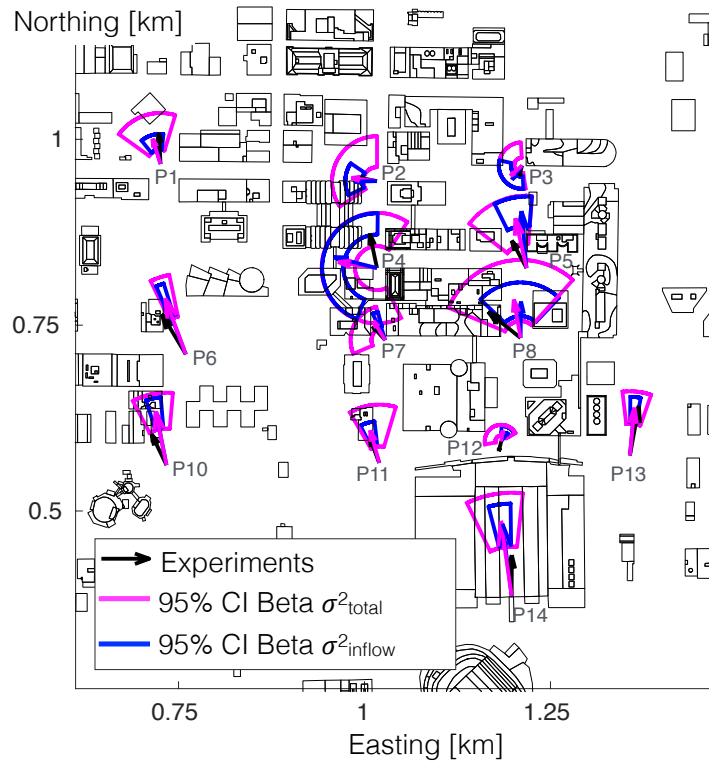


Figure 10: Comparison between experimental data and UQ study results for the wind prediction at the PWIDS stations.

434 The first observation from figure 10 is that the choice of the input pdfs
 435 for velocity magnitude and direction strongly affects the size of the 95%
 436 confidence intervals, while the prediction for the mean velocity only changes
 437 marginally. This is plausibly caused by the fact that the input pdfs differ
 438 in terms of the variance, but have the same mean value. When considering
 439 the inflow pdfs with the total variance, the 95% confidence interval from
 440 the UQ study encompasses the time-averaged measurement in 82% of the
 441 35 (PWIDS and SPWIDS) measurement stations for the wind magnitude,
 442 and in 62% of the stations for wind direction. When considering the reduced
 443 inflow variance, these numbers reduce to 65% and 41% respectively.

444 The second observation from this plot is that the UQ study clearly differ-
 445 entiates between locations in the flow with a small versus large uncertainty
 446 caused by the inflow variability. For example PWIDS 6, 10, 11 and 13 all
 447 have small 95% CIs, while other stations in the downtown area, such as

448 PWIDS 3 and 4, present much larger uncertainty induced by the inflow vari-
 449 ability. It is important to notice that this differentiation between the stations
 450 is consistent, i.e. it seems to be independent of the choice for the inflow pdfs.

451 5.2.2. Dispersion results

452 From the 10 SF_6 analysers for which the locations were shown in figure 1,
 453 only 8 were operational during the 30 minutes of interest. The mean and 95%
 454 CI for the non-dimensional concentration at these stations are presented in
 455 figure 11. The black crosses present the time-averaged measured value for K
 456 over the 30 minutes release, the magenta squares and confidence bars present
 457 the results from the UQ study using the total variance beta distributions as
 458 input, while the data plotted in blue corresponds to the inflow variance beta
 459 distribution.

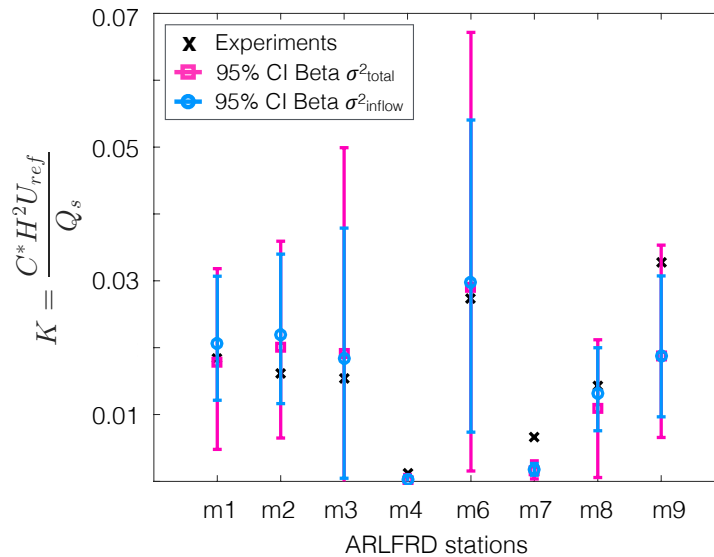


Figure 11: Comparison of the measurements during IOP9 with the uncertainty quantification study results for non dimensional concentration.

460 Similar as for the flow field, figure 11 shows that the choice of the input
 461 pdfs for velocity magnitude and direction strongly affects the size of the 95%
 462 confidence intervals, while the prediction for the mean concentration only
 463 changes marginally. Again, this is plausibly caused by the fact that the
 464 input pdfs differ in terms of the variance, but have the same mean value. In
 465 general, considering the mean from both UQ results, there is a slight over

466 prediction of the concentration in most of the stations. When using the beta
467 distributions with the total variance, the experimental mean is encompassed
468 by the 95% CI predicted by the UQ study in 75% of the stations. This number
469 decreases to 63% when using the beta distributions with the reduced inflow
470 variance.

471 The results for the concentration field also enable to differentiate between
472 locations in the flow with smaller and larger uncertainty related to variability
473 in the inflow conditions. For example, ARLFRD 1 and ARLFRD 6 exhibit
474 larger uncertainty in the predicted pollutant concentration, while ARLFRD
475 7 has a smaller confidence interval.

476 *5.2.3. Discussion on the definition of the uncertain parameter input distri-* 477 *butions*

478 The results for the flow and concentration field show a consistent be-
479 haviour in terms of the influence of the choice for the distributions for the
480 uncertain input variables. The impact on the mean prediction is negligible,
481 which can be explained by the fact that the means of the input pdfs are
482 the same. The 95% confidence intervals increase with the variance of the
483 input pdfs. As a result, the experimental mean falls within the 95% confi-
484 dence interval in more locations when using the input pdfs that represent
485 the total variance. Concluding that this approach is therefore better suited
486 for quantifying inflow uncertainties in RANS simulations of urban flow and
487 dispersion would however be premature. Instead, we argue that the input
488 pdfs with the reduced variance are a more realistic representation of the un-
489 certainty in the inflow conditions, and the resulting confidence intervals are
490 a more accurate representation of the corresponding uncertainty in the so-
491 lution. When these confidence intervals do not encompass the mean of the
492 field experiment, this is an indication that additional uncertainties are likely
493 to affect the solution. One specific example would be the uncertainty related
494 to the turbulence model. Previous research [23] has addressed quantifying
495 this uncertainty for the prediction of the flow field, and ongoing research is
496 extending this approach to the prediction of the concentration field.

497 Lastly, the results presented for the velocity and concentration field show
498 the potential of the approach to determine which areas within the city are
499 more affected by uncertainties in the inflow boundary conditions. The re-
500 sulting distinction between regions with larger and smaller uncertainties was
501 shown to be independent of the choice of the input distributions. This indi-
502 cates that the UQ approach can provide valuable information even if insuf-

503 ficient information is available to accurately characterize the input pdfs.

504 **6. Conclusions**

505 The boundary conditions for RANS simulations of urban pollutant dis-
506 persion are determined by the larger-scale atmospheric conditions, which can
507 be highly variable. The present study investigates quantifying the resulting
508 uncertainty in the prediction of urban wind and pollutant concentration. The
509 proposed method defines three uncertain input parameters that specify the
510 inflow boundary conditions: the wind velocity and direction, and the ABL
511 roughness length. A non-intrusive uncertainty quantification (UQ) approach
512 based on 729 deterministic RANS simulations is used to propagate the in-
513 flow uncertainties to the simulation output. The results of the UQ study
514 are compared to field measurements from the Joint Urban 2003 (JU2003)
515 measurement campaign.

516 An initial variance based sensitivity analysis was performed to investi-
517 gate the influence of the different uncertain input parameters on the pre-
518 diction of the QoI. The Sobol indices revealed that the prediction for the
519 non-dimensional wind speed and wind direction is primarily influenced by
520 uncertainty in the inflow wind direction, while the uncertainty in the in-
521 flow wind speed and roughness height are less influential parameters. When
522 considering tracer concentrations all three uncertain variables were shown to
523 contribute significantly to the variance in the QoI. This result points towards
524 opportunities for dimension reduction when the objective is to calculate the
525 flow field. When also considering pollutant dispersion, none of the uncer-
526 tain variables can be ignored, but the required number of evaluations could
527 potentially be reduced by using an anisotropic tensor grid [41].

528 The analysis of the results focused on comparing the prediction for the
529 mean and 95% confidence intervals to the field measurement data. We inves-
530 tigated the influence of the choice for the input pdfs for the uncertain inflow
531 velocity magnitude and direction by considering two different beta distri-
532 butions. The first represents the full variance of the experimental signal
533 measured at the most upstream location, and therefore represents all fluc-
534 tuations in the measured signal as an inflow uncertainty. The second beta
535 distribution represents a reduced variance, where the variance corresponding
536 to the turbulence kinetic energy in the RANS model is removed from the
537 measured signal, and only the remaining portion of the measured variance
538 is considered an inflow uncertainty. When using the beta distributions with

539 the total variance, the predicted 95% confidence interval encompasses the
540 time-averaged field measurement in 82% of the stations for wind magnitude,
541 62% for wind direction and 75% for the tracer concentration. These num-
542 bers respectively drop to 65%, 41%, and 63% when the beta distribution with
543 the reduced inflow variance is used. Despite this decrease in the number of
544 points where the predicted confidence intervals encompass the field measure-
545 ments, we argue that the input pdfs with the reduced variance are a more
546 realistic representation of the uncertainty in the inflow conditions. The fact
547 that the corresponding confidence intervals do not encompass the mean of
548 the field experiment, is an indication that additional uncertainties might be
549 affecting the solution. One specific example would be the uncertainty related
550 to the turbulence model, and ongoing research is focused on addressing this
551 uncertainty.

552 In summary, the results demonstrate that there is considerable uncer-
553 tainty in the prediction of urban flow and dispersion because of variability in
554 the inflow conditions. It is important to note this is an aleatory uncertainty;
555 it represents physical variability inherent to the system being analysed. It
556 is not strictly due to a lack of knowledge and can therefore not be reduced.
557 Additional experimental characterization might provide more conclusive ev-
558 idence of the variability but cannot eliminate it completely and the aleatory
559 uncertainty should therefore be characterized using probabilistic approaches.
560 The definition of the input pdfs for the uncertain variables was shown to
561 be crucial to obtain an accurate prediction for the 95% confidence interval.
562 The variance-based sensitivity analysis enabled identifying the most relevant
563 uncertain parameters, which will guide future research efforts regarding a
564 more accurate definition of the input pdfs. However, the results also show
565 that even in absence of sufficiently accurate information to define the input
566 pdfs, the proposed UQ approach can provide valuable information regarding
567 which areas in the urban canopy are more sensitive to the uncertain input
568 parameters. This information can be used to guide design decisions, such as
569 the implementation of additional geometrical features, that can reduce these
570 sensitivities and therefore establish more predictable flow and dispersion pat-
571 terns.

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- 578 [1] World Health Organization, Ambient (outdoor) air quality and health
579 (November 2016).
580 URL <http://www.who.int/mediacentre/factsheets/fs313/en/>
- 581 [2] Blocken, B., Computational fluid dynamics for urban physics: Importance, scales, possibilities, limitations and ten tips and tricks towards
582 accurate and reliable simulations, *Building and Environment* 91 (2015)
583 219–245.
584
- 585 [3] Stathopoulos, T., Pedestrian level winds and outdoor human comfort,
586 *Journal of Wind Engineering and Industrial Aerodynamics* 94 (11)
587 (2006) 769–80.
- 588 [4] Tominaga, Y., Stathopoulos, T., Turbulent Schmidt numbers for CFD
589 analysis with various types of flowfield, *Atmospheric Environment*
590 41 (37) (2007) 8091–9.
- 591 [5] Gorlé, C., van Beeck, J., Rambaud, P., Dispersion in the Wake of a Rect-
592 angular Building: Validation of Two Reynolds-Averaged Navier-Stokes
593 Modeling Approaches, *Boundary Layer Meteorology* 137 (1) (2010) 115–
594 133.
- 595 [6] Gousseau, P., Blocken, B., van Heijst, G.J.F., CFD simulation of pollu-
596 tant dispersion around isolated buildings: On the role of convective and
597 turbulent mass fluxes in the prediction accuracy, *Journal of Hazardous*
598 *Materials* 194 (2011) 422–434.
- 599 [7] Moonen, P., Gromke, C., Dorer, V., Performance assessment of Large
600 Eddy Simulation (LES) for modeling dispersion in an urban street
601 canyon with tree planting, *Atmospheric Environment* 75 (2013) 66–76.
- 602 [8] Shi, R.F., Cui, G.X., Wang, Z.S., Xu, C.X., Zhang, Z.S., Large eddy
603 simulation of wind field and plume dispersion in building array, *Atmo-
604 spheric Environment* 42 (6) (2008) 1083–1097.

- 605 [9] Boppana, V.B.L., Xie, Z.-T., Castro, I.P., Large-eddy simulation of dis-
606 persion from surface sources in arrays of obstacles, *Boundary Layer Me-*
607 *teorology* 135 (3) (2010) 433–454.
- 608 [10] Xie, Z.-T., Hayden, P. and Wood, C.R., Large-eddy simulation of
609 approaching-flow stratification on dispersion over arrays of buildings,
610 *Atmospheric Environment* 71 (2013) 64–74.
- 611 [11] Muñoz-Esparza, D., Kosovic, B., van Beeck, J., Mirocha, J.D., A
612 stochastic perturbation method to generate inflow turbulence in large-
613 eddy simulation models: Application to neutrally stratified atmospheric
614 boundary layers, *Physics of Fluids* 27 (2015) 035102.
- 615 [12] Tominaga, Y., Stathopoulos, T., CFD simulation of near-field pollu-
616 tant dispersion in the urban environment: A review of current modeling
617 techniques, *Atmospheric Environment* 79 (2013) 716–730.
- 618 [13] Hendricks, E.A., Diehl, S.R., Burrows, D.A., Keith, R., Evaluation of a
619 fast-running urban dispersion modeling system using joint urban 2003
620 field data, *Journal of Applied Meteorology and Climatology* 46 (2007)
621 2165–2179.
- 622 [14] Kumar, P., Feiz, A.-A., Ngae, P., Kumar Singh, S., Issartel, J.-P., CFD
623 simulation of short-range plume dispersion from a point release in an
624 urban like environment, *Atmospheric Environment* 122 (2015) 645–656.
- 625 [15] Chang, J.C., Hanna, S.R., Air quality model performance evaluation,
626 *Meteorology and Atmospheric Physics* 87 (1) (2004) 167–196.
- 627 [16] Emory, M., Larsson, J., Iaccarino, G., Modeling of structural uncer-
628 tainties in Reynolds-averaged Navier-Stokes closures, *Physics of Fluids*
629 25 (11) (2013) 110822.
- 630 [17] Górlé, C., Iaccarino, G., A framework for epistemic uncertainty quan-
631 tification of turbulent scalar flux models for Reynolds-averaged Navier-
632 Stokes simulations, *Physics of Fluids* 25 (2013) 055105.
- 633 [18] Górlé, C., Larsson, J., Emory, G., Iaccarino, G., The deviation from par-
634 allel shear flow as an indicator of linear eddy-viscosity model inaccuracy,
635 *Physics of Fluids* 26 (5) (2014) 051702.

- 636 [19] Xiao, H., Wu, J.-L., Wang, J.-X., Sun, R., Roy, C.J., Quantifying and
637 reducing model-form uncertainties in Reynolds-averaged Navier-Stokes
638 simulations: A data-driven, physics-informed Bayesian approach, *Journal of Computational Physics* 324 (2016) 115–136.
639
- 640 [20] Schatzmann, M., Leitl, B., Issues with validation of urban flow and
641 dispersion CFD models, *Journal of Wind Engineering and Industrial
642 Aerodynamics* 99 (2011) 169–186.
- 643 [21] Rodriguez, L.M., Bieringer, P.E., Warner, T., Urban transport and dis-
644 persion model sensitivity to wind direction uncertainty and source loca-
645 tion, *Atmospheric Environment* 64 (2013) 25–39.
- 646 [22] García-Sánchez, C., Philips, D.A., Górlé, C., Quantifying inflow uncer-
647 tainties for CFD simulations of the flow in downtown Oklahoma City,
648 *Building and Environment* 78 (118-129).
- 649 [23] Górlé, C., García-Sánchez, C., Iaccarino, G., Quantifying inflow and
650 RANS turbulence model form uncertainties for wind engineering flows,
651 *Journal of Wind Engineering and Industrial Aerodynamics* 144 (2015)
652 202–212.
- 653 [24] Allwine, K.J., Flaherty, J.E., Joint Urban 2003: Study Overview and
654 Instrument Locations, PR 15967, PNNL (2006).
- 655 [25] Chan, S.T, Leach, M.J., A validation of FEM3MP with joint urban
656 2003 data, *Journal of Applied Meteorology and Climatology* 46 (2007)
657 2127–2146.
- 658 [26] Neophytou, M., Gowardhan, A. , Brown, M., An inter-comparison of
659 three urban wind models using Oklahoma city joint urban 2003 wind
660 field measurements, *Journal of Wind Engineering and Industrial Aero-
661 dynamics* 99 (4) (2011) 357–368.
- 662 [27] Storwold, D., Meteorology division CSTE-DTC-DP-WD-ME-M, UT
663 84022-5000, Building 4034/Room 111 Dugway (2003).
- 664 [28] Wilcox, D.C., Turbulence Modeling for CFD, La Cañada: DCW Indus-
665 tries Inc., 1993.

- 666 [29] Pal Arya, S., Air pollution meteorology and dispersion, Oxford Univer-
667 sity Press, 1999.
- 668 [30] Greenshields, C.J., CFD Direct Ltd., Openfoam user guide, Tech. rep.,
669 OpenFOAM Foundation Ltd. (2011-2016).
- 670 [31] Franke, J., Best practice guideline for the CFD simulation of flows in
671 the urban environment, COST Action 732, 2007.
- 672 [32] Richards, P.J., Hoxey, R.P., Appropriate boundary conditions for com-
673 putational wind engineering models using the $k-\epsilon$ turbulence model,
674 Journal of Wind Engineering and Industrial Aerodynamics 46-47 (1993)
675 145–153.
- 676 [33] Parente, A., Górlé, C., van Beeck, J., Benocci, C., A comprehensive
677 modelling approach for the neutral atmospheric boundary layer: con-
678 sistent inflow conditions, wall function and turbulence model closure,
679 Boundary Layer Meteorology 140 (3) (2011) 411–428.
- 680 [34] Blocken, B., Stathopoulos, T., Carmeliet, J., CFD simulation of the
681 atmospheric boundary layer: wall function problems, Atmospheric En-
682 vironment 41 (2) (2007) 238–252.
- 683 [35] Burian, S.J., Han, W.S., Brown, M.J., Morphological analyses using 3D
684 building databases: Oklahoma City, Oklahoma, LANL Report LA-UR-
685 05-1821 10, Los Alamos National Laboratory (2003).
- 686 [36] Vervecken, L., Camps, J., Meyers, J., Accounting for wind-direction
687 fluctuations in reynolds-averaged simulation of near-range atmospheric
688 dispersion, Atmospheric Environment 72 (2013) 142–150.
- 689 [37] Iaccarino, G., Introduction to uncertainty representation and propa-
690 gation, AVT-193 short course on uncertainty quantification, Tech. rep.,
691 NATO Research and Technology Organization, Neuilly-sur-Seine (2011).
- 692 [38] Smith, R.C., Uncertainty Quantification: Theory, Implementation, and
693 Applications, SIAM, Computational Science and Engineering, 2014.
- 694 [39] Sobol, I.M., Sensitivity analysis for non-linear mathematical models,
695 Mathematical Modeling and Computational Experiment 1 (1993) 407–
696 414.

- 697 [40] Adams, B.M., Ebeida, M.S., Eldred, M.S., Geraci, G., Jakeman, J.D.,
698 Maupin, K.A., Monschke, J.A., Swiler, L.P., Stephens, J.A., Vigil, D.M.,
699 Wildey, T.M., Bohnhoff, W.J., Dalbey, K.R., Eddy, J.P., Hooper, R.W.,
700 Hu, K.T., Hough, P.D., Ridgwayand, E.M., Rushdi, A., Dakota, a mul-
701 tilevel parallel object-oriented framework for design optimization, pa-
702 rameter estimation, uncertainty quantification, and sensitivity analysis:
703 Version 6.3 theory manual, Tech. rep., SAND2014-4253 (2015).
- 704 [41] Paci, A., García-Sánchez, C., van Beeck, J., Gorlé, C., Uncertainty quan-
705 tification for ABL flows, Project report VKI PR 2016-20, von Karman
706 Institute for Fluid Dynamics (2015-2016).
- 707 [42] Towns, J., Cockerill, T., Dahan, M., Foster, I., Gaither, K., Grimshaw,
708 A., Hazlewood, V., Lathrop, S., Lifka, D., Peterson, G.D., Roskies, R.,
709 Scott, J.R., Wilkins-Diehr, N., XSEDE: Accelerating Scientific Discov-
710 ery, *Computing in Science & Engineering* 16 (2014) 62–74.

Highlights:

- The effect of inflow uncertainties in simulations of urban dispersion is quantified.
- The results are compared to the JU2003 experiments.
- Variance--based analysis shows wind direction is the dominant uncertain parameter.
- All uncertain parameters exhibit a non--negligible effect on urban dispersion.
- The results identify areas in the urban canopy more sensitive to inflow uncertainty.